



AFFINE TRANSLATION SURFACES WITH CONSTANT MEAN CURVATURE IN MINKOWSKI 3-SPACE

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ABSTRACT. In this paper, we obtain classification results for spacelike affine translation surfaces with constant mean curvature in three dimensional Minkowski space \mathbb{E}_1^3 .

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1. INTRODUCTION

In 3-dimensional spaces, a regular surface parameterized as $\Psi(u, v) = (u, v, z(u, v))$ is called a translation surface if usually $z(u, v)$ is of the form

$$z(u, v) = f(u) + g(v),$$

where f and g are differentiable functions of u and v , respectively. Scherk [10] discovered the first non-trivial minimal translation surface in Euclidean 3-space \mathbb{E}^3 , famously known as the Scherk surface, and is given by

$$z(u, v) = \frac{1}{c} \log \left| \frac{\cos(cu)}{\cos(cv)} \right|,$$

where $c(\neq 0)$ is a constant. Planes and Scherk surfaces are the only minimal translation surfaces in \mathbb{E}^3 . More than a century later, Liu proved that the circular cylinder is the only

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translation surface with non-zero constant mean curvature [5]. The study of constant mean curvature translation surfaces has gathered significant attention. For some of the studies, we refer the reader to see [3, 4, 5, 8, 9, 12]. A natural extension of the translation surface appears in the form of an affine translation surface, which is a surface parameterized by $\Psi(u, v) = (u, v, z(u, v))$, where now

$$z(u, v) = f(u) + g(au + v),$$

and $a(\neq 0)$ is a constant. Liu and Yu proved that the non-trivial minimal affine translation surface in \mathbb{E}^3 is given by

$$z(u, v) = \frac{1}{b} \log \left| \frac{\cos(b\sqrt{1+a^2}u)}{\cos(b(v+au))} \right|,$$

where $b(\neq 0)$ is a constant. This surface is known as the affine Scherk surface [7]. For other related works on affine translation surfaces, we refer the reader to see [1, 2, 6, 11].

In connection to the non-zero constant mean curvature of affine translation surfaces, Liu and Jung [6] obtained the classification results in \mathbb{E}^3 . Now, a Minkowski space is one of the most trivial indefinite space forms, and it marks its great significance as the trivial solution to the vacuum Einstein Field Equations without a cosmological constant. Inspired by the previous developments in the theory of constant mean curvature surfaces, we seek to classify spacelike affine translation surfaces with constant mean curvature in Minkowski 3-space \mathbb{E}_1^3 .

Consider $\Psi(u, v)$ to be a regular spacelike surface in Minkowski 3-space \mathbb{E}_1^3 . The coefficients of the 1st fundamental form E, F, G of $\Psi(u, v)$ are given by

$$E = \langle \Psi_u, \Psi_u \rangle, F = \langle \Psi_u, \Psi_v \rangle, G = \langle \Psi_v, \Psi_v \rangle,$$

and the coefficients of the 2nd fundamental form L, M, N of $\Psi(u, v)$ are given by

$$L = \langle \Psi_{uu}, \hat{n} \rangle, M = \langle \Psi_{uv}, \hat{n} \rangle, N = \langle \Psi_{vv}, \hat{n} \rangle,$$

where \hat{n} is the unit normal vector and $\langle *, * \rangle = du^2 + dv^2 - dz^2$ is the Minkowski metric. The mean curvature H of the surface $\Psi(u, v)$ is given by

$$H(u, v) = \frac{EN - 2FM + GL}{2(EG - F^2)}. \tag{1.1}$$

For a spacelike surface $\Psi(u, v)$ in \mathbb{E}_1^3 , we have $EG - F^2 > 0$, and for a timelike surfaces $EG - F^2 < 0$. In regards to the regular surface $\Psi(u, v)$ embedded in \mathbb{E}_1^3 , following two types of affine translation surfaces exist:

(i) Affine translation surface of type 1:

$$\Psi(u, v) = (u, v, z(u, v)) \quad (1.2)$$

such that

$$z(u, v) = f(u) + g(au + v). \quad (1.3)$$

(ii) Affine translation surface of type 2:

$$\Psi(u, v) = (u(v, z), v, z) \quad (1.4)$$

such that

$$u(v, z) = h(v) + t(bv + z), \quad (1.5)$$

where $a(\neq 0), b(\neq 0)$ are constants and f, g, h, t are smooth functions. We note that whenever $a = 0$ or $b = 0$, affine translation surfaces reduce simply to translation surfaces.

2. AFFINE TRANSLATION SURFACES WITH NON-ZERO CONSTANT MEAN CURVATURE

Theorem 2.1. *Let $\Psi(u, v) = (u, v, z(u, v))$ be a spacelike affine translation surface of type 1 in \mathbb{E}_1^3 . If $\Psi(u, v)$ has a non-zero constant mean curvature, then $z(u, v)$ is given by*

$$z(u, v) = \pm \frac{\sqrt{1+a^2-c^2}}{2H(1+a^2)} \sqrt{1+a^2+4H^2(b-au-v)^2} + \left(\frac{u-av}{1+a^2}\right)c + p$$

such that $c^2 < 1 + a^2$; or

$$z(u, v) = cv \pm \frac{\sqrt{1-c^2}}{2H} \sqrt{1+4H^2(b-u)^2} + q$$

such that $c^2 < 1$; where a, b, c, p, q are all constants.

Proof. We know that the mean curvature of a spacelike surface $\Psi(u, v) = (u, v, z(u, v))$ in \mathbb{E}_1^3 is given by

$$H(u, v) = \frac{(z_u^2 - 1)z_{vv} - 2z_u z_v z_{uv} + (z_v^2 - 1)z_{uu}}{2(1 - z_u^2 - z_v^2)^{\frac{3}{2}}}, \quad (2.6)$$

where z_u, z_v denotes partial differentiation of z w.r.t. u and v , respectively. We obtain the following partial derivatives of $z(u, v)$ from (1.3)

$$\begin{cases} z_u = f' + ag', \\ z_v = g', \\ z_{uu} = f'' + a^2g'', \\ z_{vv} = g'', \\ z_{uv} = ag'', \end{cases} \tag{2.7}$$

where $f' = \frac{df}{du}$ and $g' = \frac{dg}{dv}$ for $x = au + v$. Using (2.7) in (2.6), gives us

$$-f'' - (1 + a^2)g'' + (g'^2 f'' + f'^2 g'') = 2HT^3, \tag{2.8}$$

where $T^2 = 1 - (f' + ag')^2 - g'^2$ and $H(\neq 0)$ is a constant. Eqn (2.8) writes as

$$-(1 - g'^2)f'' - (1 + a^2 - f'^2)g'' = 2HT^3. \tag{2.9}$$

Now, we have the following two cases:

Case I. When $f'' = 0$, we have $f' = c$, where c is a constant. Substituting $f' = c$ in (2.9) gives us

$$-(1 + a^2 - c^2)g'' = 2H \left[1 - (c + ag')^2 - g'^2 \right]^{\frac{3}{2}}. \tag{2.10}$$

Thus, we have

$$g'' = -\frac{2H(1 + a^2)^{\frac{3}{2}}}{1 + a^2 - c^2} \left[\frac{1 + a^2 - c^2}{(1 + a^2)^2} - \left(g' + \frac{ac}{1 + a^2} \right)^2 \right]^{\frac{3}{2}}. \tag{2.11}$$

Making the following substitutions in (2.11)

$$\alpha = \frac{2H(1 + a^2)^{\frac{3}{2}}}{1 + a^2 - c^2}, \quad \beta = \frac{ac}{1 + a^2} \quad \text{and} \quad \gamma^2 = \frac{1 + a^2 - c^2}{(1 + a^2)^2},$$

results in

$$\frac{g''}{\left[\gamma^2 - (g' + \beta)^2 \right]^{\frac{3}{2}}} = -\alpha. \tag{2.12}$$

Integrating (2.12) and isolating the expression for g' gives us

$$g' = \pm \frac{\gamma^3(c_1 - \alpha x)}{\sqrt{1 + \gamma^4(c_1 - \alpha x)^2}} - \beta, \tag{2.13}$$

where c_1 is a constant. By integrating (2.13), we obtain

$$g(x) = \pm \frac{1}{\alpha\gamma} \sqrt{1 + \gamma^4(c_1 - \alpha x)^2} - \beta x + c_2, \quad (2.14)$$

where c_2 is a constant. Substituting the values of α , β and γ in (2.14) gives us

$$g(x) = \pm \frac{\sqrt{1+a^2-c^2}}{2H(1+a^2)} \sqrt{1+a^2+4H^2(c_3-x)^2} - \left(\frac{ac}{1+a^2}\right)x + c_2, \quad (2.15)$$

where c_3 is a constant. Also, $f' = c$ gives us

$$f(u) = cu + c_4, \quad (2.16)$$

where c_4 is a constant. Thus from (1.3), (2.15) and (2.16), we have

$$\begin{aligned} z(u, v) = & \pm \frac{\sqrt{1+a^2-c^2}}{2H(1+a^2)} \sqrt{1+a^2+4H^2(c_3-au-v)^2} \\ & + \left(\frac{u-av}{1+a^2}\right)c + p, \end{aligned} \quad (2.17)$$

where p is a constant and $c^2 < 1 + a^2$.

Case II. When $f'' \neq 0$. Differentiating (2.9) w.r.t. u gives us

$$\begin{aligned} (1-g'^2)f''' + (1+a^2-f'^2)ag''' - 2(f'+ag')g''f'' \\ = -6HT[(f'+ag')(f''+a^2g'') + ag'g'']. \end{aligned} \quad (2.18)$$

Now, differentiating (2.9) w.r.t. v gives us

$$(1+a^2-f'^2)g''' - 2g'g''f'' = -6HT[(f'+ag')ag'' + g'g'']. \quad (2.19)$$

Eqn's (2.18) and (2.19) yield

$$(1-g'^2)f''' - 2f'f''g'' = -6HT(f'+ag')f''. \quad (2.20)$$

Substituting the value of g'' from (2.9) in (2.20), we have

$$(1-g'^2)f''' - 2f'f'' \left[\frac{2HT^3 - (1-g'^2)}{(1+a^2-f'^2)} \right] = -6HT(f'+ag')f''. \quad (2.21)$$

Thus, we obtain

$$\begin{aligned} (1-g'^2) \left[(1+a^2-f'^2)f''' + 2f'f''^2 \right] \\ = -2HT \left[3(f'+ag')(1+a^2-f'^2) - 2T^2f' \right] f''. \end{aligned} \quad (2.22)$$

Squaring both sides of (2.22) and substituting the value of T^2 gives us

$$\begin{aligned} & (1 - g'^2)^2 \left[(1 + a^2 - f'^2) f''' + 2f' f''^2 \right]^2 \\ &= 4H^2 \left[1 - (f' + ag')^2 - g'^2 \right] \\ & \times \left[3(f' + ag')(1 + a^2 - f'^2) - 2\{1 - (f' + ag')^2 - g'^2\} f' \right]^2 f''^2. \end{aligned} \tag{2.23}$$

We notice that the above expression can be expanded as a polynomial in the powers of g' . The coefficients of g' in the above expression are functions of u , and the expression itself is identically zero, so each term must be zero. But, the coefficient of g' with the highest degree, i.e., 6 in (2.23), is $-16H^2(1 + a^2)^3 f'^2 f''^2$, which is non-zero. Thus, it follows that g' is a constant (Liu and Jung have used the same argument in [6]). Substituting $g' = c$ in (2.9) yields

$$z(u, v) = cv \pm \frac{\sqrt{1 - c^2}}{2H} \sqrt{1 + 4H^2(b - u)^2} + q, \tag{2.24}$$

where b, c, q are constants and $c^2 < 1$. Thus, the proof of the theorem is complete. □

Theorem 2.2. *Let $\Psi(v, z) = (u(v, z), v, z)$ be a spacelike affine translation surface of type 2 in \mathbb{E}_1^3 . If $\Psi(v, z)$ has a non-zero constant mean curvature, then $u(v, z)$ is given by*

$$u(v, z) = \pm \frac{\sqrt{1 - b^2 + c^2}}{2H(1 - b^2)} \sqrt{4H^2(a - bv - z)^2 - (1 - b^2)} + \left(\frac{v + bz}{1 - b^2} \right) c + p,$$

such that $1 - b^2 > 0$; or

$$u(v, z) = \pm \frac{\sqrt{1 + c^2}}{2H} \sqrt{4H^2(a - v)^2 - 1 + cz} + q,$$

where a, b, c, p, q are all constants.

Proof. The mean curvature $H(v, z)$ of a spacelike surface $r(v, z) = (u(v, z), v, z)$ in \mathbb{E}_1^3 is given by

$$H(v, z) = \frac{(u_v^2 + 1)u_{zz} - 2u_v u_z u_{vz} + (u_z^2 - 1)u_{vv}}{2(u_z^2 - u_v^2 - 1)^{\frac{3}{2}}}, \tag{2.25}$$

where u_v, u_z denotes partial differentiation of u w.r.t. v and z , respectively. We obtain the following partial derivatives of $u(v, z)$ from (1.5)

$$\begin{cases} u_v = h' + bt', \\ u_z = t', \\ u_{vv} = h'' + b^2t'', \\ u_{zz} = t'', \\ u_{vz} = bt'', \end{cases} \quad (2.26)$$

where $h' = \frac{dh}{dv}$ and $t' = \frac{dt}{dy}$ for $y = bv + z$. Using (2.26) in (2.25), gives us

$$-h'' + (1 - b^2)t'' + (t'^2h'' + h'^2t'') = 2HT^3, \quad (2.27)$$

where $T^2 = t'^2 - (h' + bt')^2 - 1$ and $H(\neq 0)$ is a constant. Eqn (2.27) writes as

$$(-1 + t'^2)h'' + (1 - b^2 + h'^2)t'' = 2HT^3, \quad (2.28)$$

Now, we have the following two cases:

Case I. When $h' = c$ is a constant. It follows from (2.28)

$$(1 - b^2 + c^2)t'' = 2H \left[t'^2 - (c + bt')^2 - 1 \right]^{\frac{3}{2}}. \quad (2.29)$$

Thus, we have

$$t'' = \frac{2H(1 - b^2)^{\frac{3}{2}}}{1 - b^2 + c^2} \left[\left(t' - \frac{bc}{1 - b^2} \right)^2 - \frac{1 - b^2 + c^2}{(1 - b^2)^2} \right]^{\frac{3}{2}}. \quad (2.30)$$

Making the following substitutions in (2.30)

$$\alpha = \frac{2H(1 - b^2)^{\frac{3}{2}}}{1 - b^2 + c^2}, \quad \beta = -\frac{bc}{1 - b^2} \quad \text{and} \quad \gamma^2 = \frac{1 - b^2 + c^2}{(1 - b^2)^2},$$

results in

$$\frac{t''}{\left[(t' + \beta)^2 - \gamma^2 \right]^{\frac{3}{2}}} = \alpha. \quad (2.31)$$

Integrating (2.31) and isolating the expression for t' gives us

$$t' = \pm \frac{\gamma^3(c_1 - \alpha y)}{\sqrt{\gamma^4(c_1 - \alpha y)^2 - 1}} - \beta, \quad (2.32)$$

where c_1 is a constant. Thereby integrating (2.32), we obtain

$$t(y) = \pm \frac{1}{\alpha\gamma} \sqrt{\gamma^4(c_1 - \alpha y)^2 - 1} - \beta y + c_2, \quad (2.33)$$

where c_2 is a constant. Substituting the values of α , β and γ in (2.33) gives us

$$t(y) = \pm \frac{\sqrt{1 - b^2 + c^2}}{2H(1 - b^2)} \sqrt{4H^2(c_3 - y)^2 - (1 - b^2)} + \frac{bc}{1 - b^2}y + c_2. \tag{2.34}$$

Also, $h' = c$ gives us

$$h(v) = cv + c_4, \tag{2.35}$$

where c_4 is a constant. Thus, from (1.5), (2.34) and (2.35), we have

$$u(v, z) = \pm \frac{\sqrt{1 - b^2 + c^2}}{2H(1 - b^2)} \sqrt{4H^2(c_3 - bv - z)^2 - (1 - b^2)} + \left(\frac{v + bz}{1 - b^2}\right) c + p, \tag{2.36}$$

where p is a constant.

Case II. When $h'' \neq 0$. Differentiating (2.28) w.r.t. v gives us

$$\begin{aligned} &(-1 + t'^2)h''' + (1 - b^2 + h'^2)bt''' + 2(h' + bt')h''t'' \\ &= 6HT [bt't'' - (h' + bt')(h'' + b^2t'')]. \end{aligned} \tag{2.37}$$

Now, differentiating (2.28) w.r.t. z gives us

$$(1 - b^2 + h'^2)t''' + 2h''t't'' = 6HT [t't'' - (h' + bt')bt'']. \tag{2.38}$$

Eqn's (2.37) and (2.38) yield

$$(-1 + t'^2)h''' + 2h'h''t'' = -6HT(h' + bt')h''. \tag{2.39}$$

Substituting the value of t'' from (2.28) in (2.39), we have

$$(-1 + t'^2)h''' + 2h'h'' \left[\frac{2HT^3 - (-1 + t'^2)}{(1 - b^2 + h'^2)} \right] = -6HT(h' + bt')h''. \tag{2.40}$$

Thus, we obtain

$$\begin{aligned} &(-1 + t'^2) \left[(1 - b^2 + h'^2)h''' - 2h'h''^2 \right] \\ &= 2HT \left[-3(h' + bt')(1 - b^2 + h'^2) - 2T^2h' \right] h'' \end{aligned} \tag{2.41}$$

Squaring both sides of (2.41) and substituting the value of T^2 gives us

$$\begin{aligned} &(-1 + t'^2)^2 \left[(1 - b^2 + h'^2)h''' - 2h'h''^2 \right]^2 \\ &= 4H^2 \left[t'^2 - (h' + bt')^2 - 1 \right] \\ &\times \left[-3(h' + bt')(1 - b^2 + h'^2) - 2\{t'^2 - (h' + bt')^2 - 1\}h' \right]^2 h''^2. \end{aligned} \tag{2.42}$$

The coefficient of t^6 in (2.42) is $16H^2(1-b^2)^3h'^2h''^2$, which is non-zero and the concluding argument in Theorem 2.1 results in t' being constant. Substituting $t' = c$ in (2.28) yields

$$u(v, z) = \pm \frac{\sqrt{1+c^2}}{2H} \sqrt{4H^2(a-v)^2 - 1 + cz + q}, \quad (2.43)$$

where a, c, q is a constant. Thus, the proof of the theorem is complete. \square

3. MAXIMAL AFFINE TRANSLATION SURFACES

Theorem 3.1. *Let $\Psi(u, v) = (u, v, z(u, v))$ be a maximal affine translation surface of type 1 in \mathbb{E}_1^3 . Then, $\Psi(u, v)$ is either a planar surface or $z(u, v)$ is given by*

$$z(u, v) = \frac{1}{c} \log \left| \frac{\cosh [c\sqrt{1+a^2}u + c_1]}{\cosh [c(au + v) + c_2]} \right| + c_3,$$

where a, c, c_1, c_2, c_3 are constants and $c \neq 0$.

Proof. Taking $H = 0$ in (2.9) gives us

$$-(1-g'^2)f'' - (1+a^2-f'^2)g'' = 0, \quad (3.44)$$

which writes as

$$\frac{f''}{1+a^2-f'^2} = \frac{-g''}{1-g'^2} = \lambda, \quad (3.45)$$

where λ is a constant and $(1+a^2-f'^2)(1-g'^2) \neq 0$.

Depending on λ , we have the following 2 cases:

Case I. $\lambda = 0$, gives us

$$f'' = 0, \quad g'' = 0, \quad (3.46)$$

which leads to a planar surface in \mathbb{E}_1^3 .

Case II. $\lambda \neq 0$, gives us

$$f(u) = \frac{1}{c} \log \left| 2 \cosh [c\sqrt{1+a^2}u + c_1] \right| + c_3, \quad (3.47)$$

$$g(au + v) = \frac{-1}{c} \log \left| 2 \cosh [c(au + v) + c_2] \right| + c_4, \quad (3.48)$$

where a, c, c_1, c_2, c_3, c_4 are constants. Thus, we have

$$z(u, v) = \frac{1}{c} \log \left| \frac{\cosh [c\sqrt{1+a^2}u + c_1]}{\cosh [c(au + v) + c_2]} \right| + p, \quad (3.49)$$

where $c \neq 0$ and p is a constant. Thus, the proof of the theorem is complete. \square

Theorem 3.2. *Let $\Psi(v, z) = (u(v, z), v, z)$ be a maximal generalized affine translation surface of type 2 in \mathbb{E}_1^3 . Then, $\Psi(v, z)$ is either a planar surface or $u(v, z)$ is given by*

$$u(v, z) = \frac{1}{c} \log \left| \frac{\cos [c\sqrt{1 - b^2}v + c_1]}{\cosh [c(bv + z) + c_2]} \right| + c_3, \quad b^2 < 1$$

or

$$u(v, z) = \frac{1}{c} \log \left| \frac{\cosh [c\sqrt{b^2 - 1}v + c_1]}{\cosh [c(bv + z) + c_2]} \right| + c_3, \quad b^2 > 1;$$

where b, c, c_1, c_2, c_3 are constants and $c \neq 0$.

Proof. For $H = 0$, it follows from (2.28)

$$(-1 + t'^2)h'' + (1 - b^2 + h'^2)t'' = 0. \tag{3.50}$$

Thus, we have

$$\frac{-h''}{1 - b^2 + h'^2} = \frac{t''}{-1 + t'^2} = \lambda, \tag{3.51}$$

where λ is a constant and $(1 - b^2 + h'^2)(-1 + t'^2) \neq 0$. $\lambda = 0$ leads to a planar surface in \mathbb{E}_1^3 and when $\lambda \neq 0$, we have the following cases:

Case I. If $b^2 < 1$; (3.51) yields

$$\begin{aligned} u(v, z) &= h(v) + t(bv + z), \\ &= \frac{1}{c} \log \left| \frac{\cos [c\sqrt{1 - b^2}v + c_1]}{\cosh [c(bv + z) + c_2]} \right| + c_3, \end{aligned} \tag{3.52}$$

where b, c, c_1, c_2, c_3 are constants and $c \neq 0$.

Case II. If $b^2 > 1$; (3.51) yields

$$\begin{aligned} u(v, z) &= h(v) + t(bv + z), \\ &= \frac{1}{c} \log \left| \frac{\cosh [c\sqrt{b^2 - 1}v + c_1]}{\cosh [c(bv + z) + c_2]} \right| + c_3, \end{aligned} \tag{3.53}$$

where b, c, c_1, c_2, c_3 are constants and $c \neq 0$. Thus, the proof of the theorem is complete. \square

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