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# A NOTE ON CSI−SUBMERSIONS FROM COSYMPLECTIC MANIFOLDS

SUSHIL KUMAR  $\mathbf{\mathbb{D}}$  \*, SUMEET KUMAR  $\mathbf{\mathbb{D}}$ , AND RAJ KUMAR SRIVASTAVA  $\mathbf{\mathbb{D}}$ 

Abstract. In this paper, our main objective is to study the notion of Clairaut semiinvariant submersions (CSI− submersions, in short) from Cosymplectic manifolds onto Riemannian manifolds. We investigate some fundamental results pertaining to the geometry of such submersions. We also obtain totally geodesicness conditions for the distributions. Moreover, we provide a non-trivial example of such Riemannian submersion.

Keywords: Riemannian submersions, Clairaut semi-invariant submersions, Almost contact metric manifolds.

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### 1. Introduction

Firstly, O' Neill [\[16\]](#page-14-0) and Gray [\[9\]](#page-14-1) separately studied the concept of Riemannian submersions between Riemannian manifolds in the 1960s. Using the notion of Riemannian submersions between almost complex manifolds, Watson [\[34\]](#page-15-0) studied almost Hermitian submersions. Further, the concept of anti-invariant submersion was first defined by Sahin [\[23\]](#page-15-1) from almost Hermitian manifolds onto Riemannian manifolds. Later, he introduced semi-invariant submersion [\[25\]](#page-15-2) from almost Hermitian manifolds onto Riemannian manifolds as a generalization of holomorphic submersions and anti-invariant submersion. Further, different kinds of Riemannian submersions on different structures have been studied, such as: slant submersions

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<sup>∗</sup> Corresponding author

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Sushil Kumar; sushilmath20@gmail.com; https://orcid.org/0000-0003-2118-4374 Sumeet Kumar; itssumeetkumar@gmail.com; https://orcid.org/0000-0003-1214-5701 Raj Kumar Srivastava; srivastavar666@gmail.com; https://orcid.org/0000-0002-2499-7402 [\[24\]](#page-15-3), semi-slant submersions [\[17\]](#page-14-2), conformal semi-slant submersion([\[12\]](#page-14-3),[\[20\]](#page-14-4)), hemi-slant Riemannian submersions [\[31\]](#page-15-4), conformal hemi-slant submersion [\[11\]](#page-14-5), quasi-bi-slant submersions [\[18\]](#page-14-6) (see also [\[13\]](#page-14-7), [\[19\]](#page-14-8), [\[21\]](#page-15-5), [\[26\]](#page-15-6), [\[28\]](#page-15-7), [\[29\]](#page-15-8)) etc.

Presently, the Riemannian submersions have abundant applications in pure mathematics and physics, for example, Kaluza-Klein theory [\[7\]](#page-14-9), Yang-Mills theory [\[8\]](#page-14-10), Supergravity and superstring theories [\[10\]](#page-14-11) etc. C. Altafini [\[2\]](#page-14-12) commenced using the notion of Riemannian submersions for the modeling and control of redundant robotic chain and proved that Riemannian submersion gives a close relationship between inverse kinematic in robotics and the pull back vectors.

In the theory of surfaces created by rotating the curves, we note that, for any geodesic  $c(c : I_1 \subset R \to \mathcal{N}_1$  on  $\mathcal{N}_1$ ) on the rotating surface  $\mathcal{N}_1$ , the product  $r \sin \Theta$  is constant along geodesic c, where  $\Theta(s)$  is the angle between  $c(s)$  and the meridian curve through  $c(s)$ ,  $s \in I_1$ , called Clairaut's theorem [\[5\]](#page-14-13). It means that it is independent of s. In 1972, Bishop [5] applied this idea to the Riemannian submersions and introduced the concept of Clairaut submersion. Afterwards, Clairaut submersions have been studied in Spacelike spaces, Timelike and Lorentzian spaces([\[15\]](#page-14-14), [\[32\]](#page-15-9), [\[33\]](#page-15-10)) and its applications in Static spacetimes [\[1\]](#page-14-15). Later on this notion has been generalized in [\[3\]](#page-14-16) and [\[15\]](#page-14-14). Kumar et al., in [\[14\]](#page-14-17), introduce the notion of Clairaut semi-invariant Riemannian map and Gupta and Singh in [\[22\]](#page-15-11) initiate the notion of Clairaut semi-invariant submersion from K¨ahler manifold and investigate some interesting geometric properties of these submersions.

In the present paper, our focus is on investigating the notion of the CSI−submersions from Cosymplectic manifolds onto Riemannian manifolds. The paper is organized as follows: In the second section, we gather main notions and formulae for other sections. In the third section, we give the definition of the CSI−submersions from Cosymplectic manifolds onto Riemannian manifolds. We investigate differential geometric properties of such submersions. In the last section, we illustrate a non-trivial example of the CSI–submersions from Cosymplectic manifolds onto Riemannian manifolds.

## 2. Preliminaries

A  $(2n + 1)$ -dimensional smooth manifold  $\mathcal{N}_1$  is said to have an almost contact structure [\[26\]](#page-15-6) if there exist on  $\mathcal{N}_1$ , a tensor field  $\phi$  of type  $(1, 1)$ , a vector field  $\xi$  and 1-form  $\eta$  such that

$$
\phi^2 = -I + \eta \otimes \xi, \ \eta \circ \phi = 0, \quad \phi \xi = 0,\tag{2.1}
$$

$$
g_1(\xi, \xi) = \eta(\xi) = 1.
$$
\n(2.2)

If there exists a Riemannian metric  $g_1$  on an almost contact manifold  $\mathcal{N}_1$  satisfying:

<span id="page-2-4"></span>
$$
g_1(\phi Z_1, \phi Z_2) = g_1(Z_1, Z_2) - \eta(Z_1)\eta(Z_2),
$$
\n
$$
g_1(Z_1, \phi Z_2) = -g_1(\phi Z_1, Z_2),
$$
\n
$$
g_1(Z_1, \xi) = \eta(Z_1),
$$
\n(2.4)

where  $Z_1, Z_2$  are any vector fields on  $\mathcal{N}_1$ . Then  $\mathcal{N}_1$  is called almost contact metric manifold [\[6\]](#page-14-18) with almost contact structure  $(\phi, \xi, \eta)$  and is represented by  $(\mathcal{N}_1, \phi, \xi, \eta, g_1)$ .

An almost contact structure  $(\phi, \xi, \eta)$  is said to be normal if the almost complex structure J on the product manifold  $\mathcal{N}_1 \times R$  is given by

$$
J(Z_1, \mathcal{F}\frac{d}{dt}) = (\phi Z_1 - \mathcal{F}\xi, \eta(Z_1)\frac{d}{dt}),\tag{2.5}
$$

where  $J^2 = -I$  and F is a differentiable function on  $\mathcal{N}_1 \times R$  that has no torsion, i.e., J is integrable. The form for normality in terms of  $\phi$ ,  $\xi$  and  $\eta$  is given by  $[\phi, \phi] + 2d\eta \otimes \xi = 0$  on  $\mathcal{N}_1$ , where  $[\phi, \phi]$  is the Nijenhuis tensor of  $\phi$ . Further, the fundamental 2-form  $\Phi$  is defined by  $\Phi(Z_1, Z_2) = g_1(Z_1, \phi Z_2)$ .

A manifold  $\mathcal{N}_1$  with the structure  $(\phi, \xi, \eta, g_1)$  is said to be Cosymplectic [\[26\]](#page-15-6) if

<span id="page-2-3"></span>
$$
(\nabla_{Z_1}\phi)Z_2 = 0\tag{2.6}
$$

for any vector fields  $Z_1, Z_2$  on  $\mathcal{N}_1$ , where  $\nabla$  stands for the Riemannian connection of the metric  $g_1$  on  $\mathcal{N}_1$ . For a Cosymplectic manifold, we have

<span id="page-2-2"></span>
$$
\nabla_{Z_1}\xi = 0\tag{2.7}
$$

for any vector field  $Z_1$  on  $\mathcal{N}_1$ .

O'Neill's tensors [\[16\]](#page-14-0)  $\mathcal T$  and  $\mathcal A$  are given by

<span id="page-2-0"></span>
$$
\mathcal{A}_{X_1} X_2 = \mathcal{H} \nabla_{\mathcal{H} X_1} \mathcal{V} X_2 + \mathcal{V} \nabla_{\mathcal{H} X_1} \mathcal{H} X_2, \tag{2.8}
$$

<span id="page-2-1"></span>
$$
\mathcal{T}_{X_1} X_2 = \mathcal{H} \nabla_{\mathcal{V} X_1} \mathcal{V} X_2 + \mathcal{V} \nabla_{\mathcal{V} X_1} \mathcal{H} X_2 \tag{2.9}
$$

for any  $X_1, X_2$  on  $\mathcal{N}_1$ . For vertical vector fields  $Y_1, Y_2$ , the tensor field  $\mathcal T$  has the symmetry property, that is,

$$
\mathcal{T}_{Y_1} Y_2 = \mathcal{T}_{Y_2} Y_1,\tag{2.10}
$$

while for horizontal vector fields  $X_1, X_2$ , the tensor field  $\mathcal A$  has alternation property, that is

$$
A_{X_1} X_2 = -A_{X_2} X_1. \t\t(2.11)
$$

From the equations  $(2.8)$  and  $(2.9)$  $(2.9)$ , we have

<span id="page-3-0"></span>
$$
\nabla_{Y_1} Y_2 = \mathcal{T}_{Y_1} Y_2 + \mathcal{V} \nabla_{Y_1} Y_2, \tag{2.12}
$$

<span id="page-3-5"></span>
$$
\nabla_{Y_1} Z_1 = \mathcal{T}_{Y_1} Z_1 + \mathcal{H} \nabla_{Y_1} Z_1, \tag{2.13}
$$

<span id="page-3-3"></span>
$$
\nabla_{Z_1} Y_1 = \mathcal{A}_{Z_1} Y_1 + \mathcal{V} \nabla_{Z_1} Y_1, \tag{2.14}
$$

<span id="page-3-1"></span>
$$
\nabla_{Z_1} Z_2 = \mathcal{H} \nabla_{Z_1} Z_2 + \mathcal{A}_{Z_1} Z_2 \tag{2.15}
$$

for all  $Y_1, Y_2 \in \Gamma(\ker F_*)$  and  $Z_1, Z_2 \in \Gamma(\ker F_*)^{\perp}$ , where  $\mathcal{H} \nabla_{Y_1} Z_1 = \mathcal{A}_{Z_1} Y_1$ , if  $Z_1$  is basic. It can be easily seen that  $\mathcal T$  acts on the fibers as the second fundamental form, while  $\mathcal A$ acts on the horizontal distribution and measures the obstruction to the integrability of this distribution.

The Riemannian submersion  $F$  between two Riemannian manifolds is totally geodesic if

$$
(\nabla F_*)(U_1, U_2) = 0 \quad \forall \quad U_1, U_2 \in \Gamma(T\mathcal{N}_1).
$$

Totally umbilical Riemannian submersion is a Riemannian submersion with totally umbilicalfibers  $([4], [5])$  $([4], [5])$  $([4], [5])$  $([4], [5])$  $([4], [5])$  if

<span id="page-3-2"></span>
$$
\mathcal{T}_{Z_1} Z_2 = g_1(Z_1, Z_2) H \tag{2.16}
$$

for all  $Z_1, Z_2 \in \Gamma(\ker F_*)$ , where H denotes the mean curvature vector field of fibers.

Let  $F: (\mathcal{N}_1, g_1) \to (\mathcal{N}_2, g_2)$  be a Riemannian submersion between Riemannian manifolds. The differential map  $F_*$  of F can be viewed as a section of the bundle  $Hom(T\mathcal{N}_1, F^{-1}T\mathcal{N}_2) \rightarrow$  $\mathcal{N}_1$ , where  $F^{-1}T\mathcal{N}_2$  is the pullback bundle whose fibers at  $q\in \mathcal{N}_1$  is  $(F^{-1}T\mathcal{N}_2)_q=T_F(q)\mathcal{N}_2,q\in \mathcal{N}_1$  $\mathcal{N}_1$ . The bundle  $Hom(T\mathcal{N}_1, F^{-1}T\mathcal{N}_2)$  has a connection  $\nabla$  induced from the Levi-Civita connection  $\nabla^{\mathcal{N}_1}$  and the pullback connection  $\nabla^F$ . Then the second fundamental form of F is given by

<span id="page-3-4"></span>
$$
(\nabla F_*)(V_1, V_2) = \nabla_{V_1}^F F_*(V_2) - F_*(\nabla_{V_1}^{V_1} V_2)
$$
\n(2.17)

for the vector fields  $V_1, V_2 \in \Gamma(T\mathcal{N}_1)$ .

### 3. The CSI−submersions from Cosymplectic manifolds

In the theory of Riemannian submersions, Bishop [\[5\]](#page-14-13) initiated the concept of Clairaut submersion as: a submersion  $F : (\mathcal{N}_1, g_1) \to (\mathcal{N}_2, g_2)$  is called a Clairaut submersion if there exist a function  $r : \mathcal{N}_1 \to R^+$  in such a way that any geodesic that makes an angle  $\Theta$  with a horizontal subspace,  $r \sin \Theta$  is constant.

On the other side, Sahin [\[27\]](#page-15-12) generalized the concept of Clairaut submersion and initiated the study of Clairaut Riemannian maps and investigated its geometric properties.

**Theorem 3.1.** [\[5\]](#page-14-13) Let  $F : (\mathcal{N}_1, g_1) \to (\mathcal{N}_2, g_2)$  be a Riemannian submersion with connected fibers. Then, F is a Clairaut Riemannian submersion with  $r = e^h$  if each fiber is totally umbilical and has the mean curvature vector field  $H = -\nabla h$ , where  $\nabla h$  is the gradient of the function h with respect to  $q_1$ .

**Definition 3.1.** [\[26\]](#page-15-6) Let F be a Riemannian submersion from an almost contact metric manifold  $(\mathcal{N}_1, \phi, \xi, \eta, g_1)$  onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$ . Then we say that F is a semi-invariant submersion if there is a distribution  $\mathfrak{D}_1 \subseteq \ker F_*$  such that

$$
\ker F_* = \mathfrak{D}_1 \oplus \mathfrak{D}_2, \quad \phi(\mathfrak{D}_1) = \mathfrak{D}_1, \phi(\mathfrak{D}_2) \subseteq (\ker F_*)^{\perp},
$$

where  $\mathfrak{D}_1$  and  $\mathfrak{D}_2$  mutually orthogonal distributions in (ker  $F_*$ ).

Let  $\mu$  denotes the complementary orthogonal subbundle to  $\phi(\mathfrak{D}_2)$  in  $(\ker F_*)^{\perp}$ . Then we have

$$
(\ker F_*)^{\perp} = \phi(\mathfrak{D}_2) \oplus \mu.
$$

Obviously  $\mu$  is an invariant subbundle of  $(\ker F_*)^{\perp}$  with respect to the contact structure  $\phi$ .

We say that a semi-invariant submersion  $F : \mathcal{N}_1 \to \mathcal{N}_2$  admits a vertical Reeb vector field  $\xi$  if it is tangent to (ker  $F_*$ ) and it admits horizontal Reeb vector field  $\xi$  it is normal to (ker  $F_$ ). One can easily observe that  $\mu$  contains the Reeb vector field in case of the Riemannian submersion admits horizontal Reeb vector field.

We now define the notion of CSI− submersion in contact manifolds as follows:

**Definition 3.2.** A semi-invariant submersions from a Cosymplectic manifold  $(N_1, \phi, \xi, \eta, g_1)$ onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$  is called CSI– submersion if it satisfies the condition of Clairaut Riemannian submersion.

For any vector field  $W_1 \in \Gamma(\ker F_*)$ , we put

$$
W_1 = PW_1 + QW_1, \t\t(3.18)
$$

where P and Q are projection morphisms [\[4\]](#page-14-19) of ker  $F_*$  onto  $\mathfrak{D}_1$  and  $\mathfrak{D}_2$ , respectively.

For  $U_1 \in (\ker F_*)$ , we get

<span id="page-4-0"></span>
$$
\phi U_1 = \psi U_1 + \omega U_1,\tag{3.19}
$$

where  $\psi U_1 \in \Gamma(\mathfrak{D}_1)$  and  $\omega U_1 \in \Gamma(\phi \mathfrak{D}_2)$ . Also for  $V_2 \in \Gamma(\ker F_*)^{\perp}$ , we get

<span id="page-4-1"></span>
$$
\phi V_2 = BV_2 + CV_2,\tag{3.20}
$$

where  $BV_2 \in \Gamma(\mathfrak{D}_2)$  and  $CV_2 \in \Gamma(\mu)$ .

Definition 3.3. [\[30\]](#page-15-13) Let F be a CSI− submersion from an almost contact metric manifold  $(\mathcal{N}_1, \phi, \xi, \eta, g_1)$  onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$ . If  $\mu = \{0\}$  or  $\mu = \xi >$ , i.e.,  $(\ker F_*)^{\perp} = \phi(\mathfrak{D}_2)$  or  $(\ker F_*)^{\perp} = \phi(\mathfrak{D}_2) \oplus \langle \xi \rangle$  respectively, then we call  $\phi$  a Lagrangian Riemannian submersion. In this case, for any horizontal vector field  $Z_1$ , we have

$$
BZ_1 = \phi Z_1 \text{ and } CZ_1 = 0. \tag{3.21}
$$

**Lemma 3.1.** Let F be a CSI- submersion from a Cosymplectic manifold  $(N_1, \phi, \xi, \eta, g_1)$ onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$  admitting vertical or horizontal Reeb vector field. Then, we get

$$
\mathcal{V}\nabla_{W_1}\psi W_2 + \mathcal{T}_{W_1}\omega W_2 = B\mathcal{T}_{W_1}W_2 + \psi \mathcal{V}\nabla_{W_1}W_2, \tag{3.22}
$$

$$
\mathcal{T}_{W_1}\psi W_2 + \mathcal{H}\nabla_{W_1}\omega W_2 = C\mathcal{T}_{W_1}W_2 + \omega\mathcal{V}\nabla_{W_1}W_2, \tag{3.23}
$$

$$
\mathcal{V}\nabla_{U_1}BU_2 + \mathcal{A}_{U_1}CU_2 = B\mathcal{H}\nabla_{U_1}U_2 + \psi \mathcal{A}_{U_1}U_2, \tag{3.24}
$$

$$
\mathcal{A}_{U_1}BU_2 + \mathcal{H}\nabla_{U_1}CU_2 = C\mathcal{H}\nabla_{U_1}U_2 + \omega \mathcal{A}_{U_1}U_2,
$$
\n(3.25)

$$
\mathcal{V}\nabla_{W_1}BU_1 + \mathcal{T}_{W_1}CU_1 = \psi \mathcal{T}_{W_1}U_1 + B\mathcal{H}\nabla_{W_1}U_1, \tag{3.26}
$$

$$
\mathcal{T}_{W_1}BU_1 + \mathcal{H}\nabla_{W_1}CU_1 = \omega \mathcal{T}_{W_1}U_1 + C\mathcal{H}\nabla_{W_1}U_1, \tag{3.27}
$$

$$
\mathcal{V}\nabla_{U_1}\psi W_1 + \mathcal{A}_{U_1}\omega W_1 = B\mathcal{A}_{U_1}W_1 + \psi \mathcal{V}\nabla_{U_1}W_1, \tag{3.28}
$$

$$
\mathcal{A}_{U_1}\psi W_1 + \mathcal{H}\nabla_{U_1}\omega W_1 = C\mathcal{A}_{U_1}W_1 + \omega\mathcal{V}\nabla_{U_1}W_1, \tag{3.29}
$$

where  $W_1, W_2 \in \Gamma(\ker F_*)$  and  $U_1, U_2 \in \Gamma(\ker F_*)^{\perp}$ .

**Proof.** Using  $(2.12) - (2.15)$  $(2.12) - (2.15)$ ,  $(3.19)$  and  $(3.20)$ , we get Lemma 3.1.

**Corollary 3.1.** Let F be a Lagrangian submersion from a Cosymplectic manifold  $(\mathcal{N}_1, \phi, \xi, \eta, \phi)$  $g_1$ ) onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$  admitting vertical or horizontal Reeb vector field. Then we get

$$
\mathcal{V}\nabla_{V_1}\psi V_2 + \mathcal{T}_{V_1}\omega V_2 = B\mathcal{T}_{V_1}V_2 + \psi\mathcal{V}\nabla_{V_1}V_2, \mathcal{T}_{V_1}\psi V_2 + \mathcal{H}\nabla_{V_1}\omega V_2 = \omega\mathcal{V}\nabla_{V_1}V_2,
$$
  
\n
$$
\mathcal{V}\nabla_{Y_1}BY_2 = B\mathcal{H}\nabla_{Y_1}Y_2 + \psi\mathcal{A}_{Y_1}Y_2, \mathcal{A}_{Y_1}BY_2 = \omega\mathcal{A}_{Y_1}Y_2,
$$
  
\n
$$
\mathcal{V}\nabla_{V_1}BY_1 = \psi\mathcal{T}_{V_1}Y_1 + B\mathcal{H}\nabla_{V_1}Y_1, \mathcal{T}_{V_1}BY_1 = \omega\mathcal{T}_{V_1}Y_1,
$$
  
\n
$$
\mathcal{V}\nabla_{Y_1}\psi V_1 + \mathcal{A}_{Y_1}\omega V_1 = B\mathcal{A}_{Y_1}V_1 + \psi\mathcal{V}\nabla_{Y_1}V_1, \mathcal{A}_{Y_1}\psi V_1 + \mathcal{H}\nabla_{Y_1}\omega V_1 = \omega\mathcal{V}\nabla_{Y_1}V_1,
$$
  
\nwhere  $V_1, V_2 \in \Gamma(\ker F_*)$  and  $Y_1, Y_2 \in \Gamma(\ker F_*)^{\perp}.$ 

**Lemma 3.2.** Let F be a CSI- submersion from a Cosymplectic manifold  $(N_1, \phi, \xi, \eta, g_1)$ onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$  admitting vertical or horizontal Reeb vector field. Then we have

$$
\mathcal{T}_{Z_1}\xi = 0, \mathcal{A}_{Z_2}\xi = 0 \tag{3.30}
$$

for  $Z_1 \in \Gamma(\ker F_*)^{\perp}$  and  $Z_2 \in \Gamma(\ker F_*)^{\perp}$ .

**Proof.** Using  $(2.12)–(2.15)$  $(2.12)–(2.15)$  and  $(2.7)$ , we get Lemma 3.2.

**Lemma 3.3.** Let F be a CSI- submersion from a Cosymplectic manifold  $(N_1, \phi, \xi, \eta, g_1)$ onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$ . If  $\gamma : I_2 \subset R \to \mathcal{N}_1$  is a regular curve and  $Z_1(t)$  and  $Z_2(t)$  are the vertical and horizontal components of the tangent vector field  $\dot{\gamma} = E$  of  $\gamma(t)$ , respectively, then  $\gamma$  is a geodesic if and only if along  $\gamma$  the following equations hold:

<span id="page-6-0"></span>
$$
\mathcal{V}\nabla_{\dot{\gamma}}BZ_2 + \mathcal{V}\nabla_{\dot{\gamma}}\psi Z_1 + (\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})CZ_2 + (\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})\omega Z_1 = 0, \qquad (3.31)
$$

<span id="page-6-1"></span>
$$
\mathcal{H}\nabla_{\dot{\gamma}}CZ_2 + \mathcal{H}\nabla_{\dot{\gamma}}\omega Z_1 + (\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})BZ_2 + (\mathcal{A}_{Z_2} + \mathcal{T}_{Z_1})\psi Z_1 = 0.
$$
 (3.32)

**Proof.** Let  $\gamma: I_2 \to \mathcal{N}_1$  be a regular curve on  $\mathcal{N}_1$ . Since  $\dot{\gamma}(t) = Z_1(t) + Z_2(t)$ , where  $Z_1(t)$  and  $Z_2(t)$  are the vertical and horizontal components of  $\dot{\gamma}(t)$ . Using  $(2.6),(2.12)-(2.15)$  $(2.6),(2.12)-(2.15)$  $(2.6),(2.12)-(2.15)$  $(2.6),(2.12)-(2.15)$ , (3.[19\)](#page-4-0) and [\(3.20\)](#page-4-1), we have

$$
\phi \nabla_{\dot{\gamma}} \dot{\gamma} = \nabla_{\dot{\gamma}} \phi \dot{\gamma}
$$
  
\n
$$
= \nabla_{Z_1} \psi Z_1 + \nabla_{Z_1} \omega Z_1 + \nabla_{Z_2} \psi Z_1 + \nabla_{Z_2} \omega Z_1 +
$$
  
\n
$$
\nabla_{Z_1} B Z_2 + \nabla_{Z_1} C Z_2 + \nabla_{Z_2} B Z_2 + \nabla_{Z_2} C Z_2,
$$
  
\n
$$
= \mathcal{V} \nabla_{\dot{\gamma}} B Z_2 + \mathcal{V} \nabla_{\dot{\gamma}} \psi Z_1 + (\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2}) C Z_2 + (\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2}) \omega Z_1 +
$$
  
\n
$$
\mathcal{H} \nabla_{\dot{\gamma}} C Z_2 + \mathcal{H} \nabla_{\dot{\gamma}} \omega Z_1 + (\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2}) B Z_2 + (\mathcal{A}_{Z_2} + \mathcal{T}_{Z_1}) \psi Z_1.
$$

From above, vertical and horizontal components are:

$$
\mathcal{V}\phi\nabla_{\dot{\gamma}}\dot{\gamma} = \mathcal{V}\nabla_{\dot{\gamma}}BZ_2 + \mathcal{V}\nabla_{\dot{\gamma}}\psi Z_1 + (\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})CZ_2 + (\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})\omega Z_1,
$$
  

$$
\mathcal{H}\phi\nabla_{\dot{\gamma}}\dot{\gamma} = \mathcal{H}\nabla_{\dot{\gamma}}CZ_2 + \mathcal{H}\nabla_{\dot{\gamma}}\omega Z_1 + (\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})BZ_2 + (\mathcal{A}_{Z_2} + \mathcal{T}_{Z_1})\psi Z_1.
$$

Thus  $\gamma$  is a geodesic on  $\mathcal{N}_1$  if and only if  $\mathcal{V}\phi\nabla_{\dot{\gamma}}$  $\dot{\gamma} = 0$  and  $\mathcal{H} \phi \nabla_{\dot{\gamma}}$  $\dot{\gamma}=0.$ 

**Theorem 3.2.** Let  $F$  be a Clairaut semi-invariant submersion from a Cosymplectic manifold  $(N_1, \phi, \xi, \eta, g_1)$  onto a Riemannian manifold  $(N_2, g_2)$ . Then F is a CSI- submersion with  $r = e^h$  if and only if

$$
g_1(\nabla h, Z_2)||Z_1||^2 = g_1(\mathcal{V}\nabla_{\dot{\gamma}}BZ_2, \psi Z_1) + g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})CZ_2, \psi Z_1) + g_1(\mathcal{H}\nabla_{\dot{\gamma}}CZ_2, \omega Z_1) + g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})BZ_2, \omega Z_1),
$$

where  $\gamma: I_2 \to \mathcal{N}_1$  is a geodesic on  $\mathcal{N}_1$ ,  $Z_1(t)$  and  $Z_2(t)$  are vertical and horizontal components of  $\dot{\gamma}(t)$ , respectively.

**Proof.** Let  $\gamma: I_2 \to \mathcal{N}_1$  be a geodesic on  $\mathcal{N}_1$  with  $Z_1(t) = \mathcal{V}\dot{\gamma}(t)$  and  $Z_2(t) = \mathcal{H}\dot{\gamma}(t)$ . Let  $\Theta(t)$  denotes the angle in  $[0, \pi]$  between  $\dot{\gamma}(t)$  and  $Z_2(t)$ . Assuming  $v = ||\dot{\gamma}(t)||^2$ , then we get

<span id="page-7-0"></span>
$$
g_1(Z_1(t), Z_1(t)) = v \sin^2 \Theta(t), \tag{3.33}
$$

$$
g_1(Z_2(t), Z_2(t)) = v \cos^2 \Theta(t). \tag{3.34}
$$

Now, differentiating [\(3.33\)](#page-7-0), we get

$$
\frac{d}{dt}g_1(Z_1(t), Z_1(t)) = 2v \sin \Theta(t) \cos \Theta(t) \frac{d\Theta}{dt},
$$
  

$$
g_1(\nabla_\gamma Z_1(t), Z_1(t)) = v \cos \Theta(t) \cos \Theta(t) \frac{d\Theta}{dt}.
$$

Using equations  $(2.3)$  and  $(2.6)$  in above equation, we get

<span id="page-7-2"></span>
$$
g_1(\nabla_{\dot{\gamma}} \phi Z_1(t), \phi Z_1(t)) = \upsilon \sin \Theta(t) \cos \Theta(t) \frac{d\Theta}{dt}.
$$
\n(3.35)

Now we obtain

<span id="page-7-1"></span>
$$
g_1(\nabla_{\dot{\gamma}} \phi Z_1, \phi Z_1) = g_1(\mathcal{V} \nabla_{\dot{\gamma}} \psi Z_1, \psi Z_1) + g_1(\mathcal{H} \nabla_{\dot{\gamma}} \omega Z_1, \omega Z_1) + g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2}) \psi Z_1, \omega Z_1) + g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2}) \omega Z_1, \psi Z_1).
$$
\n(3.36)

Using equations  $(3.31)$  and  $(3.32)$  in  $(3.37)$ , we have

<span id="page-7-3"></span>
$$
g_1(\nabla_{\dot{\gamma}} \phi Z_1, \phi Z_1) = -g_1(\mathcal{V} \nabla_{\dot{\gamma}} B Z_2, \psi Z_1) - g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2}) C Z_2, \psi Z_1) - (3.37)
$$
  

$$
g_1(\mathcal{H} \nabla_{\dot{\gamma}} C Z_2, \omega Z_1) - g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2}) B Z_2, \omega Z_1).
$$

From  $(3.35)$  and  $(3.38)$ , we have

<span id="page-7-4"></span>
$$
v \cos \Theta(t) \cos \Theta(t) \frac{d\Theta}{dt} = -g_1(\mathcal{V} \nabla_{\dot{\gamma}} B Z_2, \psi Z_1) - g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2}) C Z_2, \psi Z_1) - (3.38)
$$

$$
g_1(\mathcal{H} \nabla_{\dot{\gamma}} C Z_2, \omega Z_1) - g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2}) B Z_2, \omega Z_1).
$$

Moreover,  $\pi$  is a CSI– Riemannian submersion with  $r = e^h$  if and only if  $\frac{d}{dt}(e^{h \circ \gamma} \sin \Theta) =$ 0, i.e.,  $e^{h\circ\gamma}(\cos\Theta \frac{d\Theta}{dt} + \sin\Theta \frac{dh}{dt}) = 0$ . By multiplying this with non-zero factor  $v \sin \Theta$ , we have

<span id="page-8-0"></span>
$$
-v \cos \Theta \sin \Theta \frac{d\Theta}{dt} = v \sin^2 \Theta \frac{dh}{dt},
$$
  
\n
$$
v \cos \Theta \sin \Theta \frac{d\Theta}{dt} = -g_1(Z_1, Z_1) \frac{dh}{dt},
$$
  
\n
$$
v \cos \Theta \sin \Theta \frac{d\Theta}{dt} = -g_1(\nabla h, \dot{\gamma}) ||Z_1||^2,
$$
  
\n
$$
v \cos \Theta \sin \Theta \frac{d\Theta}{dt} = -g_1(\nabla h, Z_2) ||Y_1||^2.
$$
\n(3.39)

Thus, from equations [\(3.39\)](#page-7-4) and [\(3.39\)](#page-8-0), we have

$$
g_1(\nabla h, Z_2)||Z_1||^2 = g_1(\mathcal{V}\nabla_{\dot{\gamma}}BZ_2, \psi Z_1) + g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})CZ_2, \psi Z_1) + g_1(\mathcal{H}\nabla_{\dot{\gamma}}CZ_2, \omega Z_1) + g_1((\mathcal{T}_{Z_1} + \mathcal{A}_{Z_2})BZ_2, \omega Z_1).
$$

Hence Theorem 3.2 is proved.

**Corollary 3.2.** Let F be a CSI- submersion from a Cosymplectic manifold  $(\mathcal{N}_1, \phi, \xi, \eta, g_1)$ to a Riemannian manifold  $(\mathcal{N}_2, g_2)$  admitting horizontal Reeb vector field. Then we get

$$
g_1(\nabla h, \xi) = 0.
$$

**Theorem 3.3.** Let F be a CSI- submersion from a Cosymplectic manifold  $(N_1, \phi, \xi, \eta, g_1)$ onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$  with  $r = e^h$ , then at least one of the following statement is true:

- (i) h is constant on  $\phi(\mathfrak{D}_2)$ ,
- $(ii)$  the fibers are one-dimensional,

$$
(iii) \nabla_{\phi X_1} F_*(W_1) = -W_1(h) F_*(\phi X_1), \text{ for all } X_1 \in \Gamma(\mathfrak{D}_2), W_1 \in \Gamma(\mu) \text{ and } \xi \neq W_1.
$$

**Proof.** Let F be  $CSI$  – submersion from a Cosymplectic manifold onto a Riemannian manifold. For  $Y_1, Y_2 \in \Gamma(\mathfrak{D}_2)$ , using [\(2.16\)](#page-3-2) and Theorem 3.1, we get

<span id="page-8-1"></span>
$$
\mathcal{T}_{Y_1} Y_2 = -g_1(Y_1, Y_2) \text{grad} h. \tag{3.40}
$$

Taking the inner product in  $(3.40)$  with  $\phi X_1$ , we get

<span id="page-8-2"></span>
$$
g_1(\mathcal{T}_{Y_1} Y_2, \phi X_1) = -g_1(Y_1, Y_2)g_1(gradh, \phi X_1)
$$
\n(3.41)

for all  $X_1 \in \Gamma(\mathfrak{D}_2)$ .

From  $(2.3), (2.6)$  $(2.3), (2.6)$  $(2.3), (2.6)$  $(2.3), (2.6)$  and  $(3.41)$  $(3.41)$ , we obtain

$$
g_1(\nabla_{Y_1} \phi Y_2, X_1) = g_1(Y_1, Y_2)g_1(gradh, \phi X_1).
$$

By using  $(2.3)$  $(2.3)$  and  $(2.16)$  $(2.16)$  in above equation, we have

<span id="page-9-0"></span>
$$
g_1(Y_1, X_1)g_1(gradh, \phi Y_2) = g_1(Y_1, Y_2)g_1(gradh, \phi X_1).
$$
\n(3.42)

Taking  $X_1 = Y_2$  and interchanging the role of  $Y_1$  and  $Y_2$ , we get

<span id="page-9-1"></span>
$$
g_1(Y_2, Y_2)g_1(gradh, \phi Y_1) = g_1(Y_1, Y_2)g_1(gradh, \phi Y_2). \tag{3.43}
$$

Using (3.[42\)](#page-9-0) with  $X_1 = Y_1$  in (3.[43\)](#page-9-1), we have

<span id="page-9-2"></span>
$$
g_1(gradh, \phi Y_1) = \frac{(g_1(Y_1, Y_2))^2}{||Y_1||^2||Y_2||^2} g_1(gradh, \phi Y_1).
$$
\n(3.44)

If  $gradh \in \Gamma(\phi(\mathfrak{D}_2))$ , then (3.[44\)](#page-9-2) and the equality condition of Schwarz inequality implies that either h is constant on  $\phi(\mathfrak{D}_2)$  or the fibers are 1-dimensional. This implies the proof of  $(i)$  and  $(ii)$ .

Now, from (2.[15\)](#page-3-1) and (2.[16\)](#page-3-2), we get

<span id="page-9-3"></span>
$$
g_1(\nabla_{Y_1} X_1, W_1) = -g_1(Y_1, X_1)g_1(gradh, W_1),
$$
\n(3.45)

for all  $W_1 \in \Gamma(\mu)$  and  $\xi \neq W_1$ . Using  $(2.3), (2.6)$  $(2.3), (2.6)$  $(2.3), (2.6)$  $(2.3), (2.6)$  $(2.3), (2.6)$  and  $(3.45)$  $(3.45)$ , we have

$$
g_1(\nabla_{Y_1}\phi X_1, \phi W_1) = -g_1(Y_1, X_1)g_1(gradh, W_1),
$$

which implies

<span id="page-9-4"></span>
$$
g_1(\nabla_{\phi X_1} Y_1, \phi W_1) = -g_1(Y_1, X_1)g_1(gradh, W_1).
$$
\n(3.46)

By using  $(2.14)$  $(2.14)$  and  $(3.46)$  $(3.46)$ , we have

$$
g_1(\mathcal{H}\nabla_{\phi X_1} W_1, \phi Y_1) = -g_1(\phi Y_1, \phi X_1)g_1(\text{ grad } h, W_1).
$$

Also for Riemannian submersion  $F$ , we have

<span id="page-9-5"></span>
$$
g_2(F_*(\nabla_{\phi X_1}^{X_1} W_1), F_*(\phi Y_1)) = -g_2(F_*(\phi Y_1), F_*(\phi X_1))g_1(gradh, W_1). \tag{3.47}
$$

Again, using  $(2.17)$  $(2.17)$  and  $(3.47)$  $(3.47)$ , we get

$$
g_2(\nabla_{\phi X_1} F_*(W_1), F_*(\phi Y_1)) = -g_2(F_*(\phi Y_1), F_*(\phi X_1))g_1(gradh, W_1),
$$

which implies.

<span id="page-9-6"></span>
$$
\nabla_{\phi X_1} F_*(W_1) = -W_1(h) F_*(\phi X_1). \tag{3.48}
$$

If gradh  $\in \Gamma(\mu) \backslash {\xi}$ , then (3.[48\)](#page-9-6) implies (iii).

Corollary 3.3. Let F be a CSI – submersion from a Cosymplectic manifold  $(N_1, \phi, \xi, \eta, g_1)$ onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$  with  $r = e^h$  and  $\dim(D_2) > 1$ . Then the fibers of F are totally geodesic if and only if  $\overline{\nabla}_{\phi X_1} F_*(W_1) = 0 \ \forall X_1 \in \Gamma(\mathfrak{D}_2)$  and  $W_1 \in \Gamma(\mu)$ .

**Lemma 3.4.** Let F be a CSI- submersion from a Cosymplectic manifold  $(N_1, \phi, \xi, \eta, g_1)$ onto a Riemannian manifold  $(N_2, g_2)$  with  $r = e^h$  and  $\dim(D_2) > 1$ . Then  $\overline{\nabla}_{W_1} F_*(\phi Y_1) =$  $W_1(h)F_*(\phi Y_1)$  for  $Y_1 \in \Gamma(\mathfrak{D}_2)$  and  $W_1 \in \Gamma(\ker F_*)^{\perp} \setminus \{\xi\}.$ 

**Proof.** Let F be a  $CSI$ - submersion from a Cosymplectic manifold onto a Riemannian manifold. From Theorem 3.1, fibers are totally umbilical with mean curvature vector field  $H = -gradh$ , then we get

$$
-g_1(\nabla_{Y_1} W_1, Y_2) = g_1(\nabla_{Y_1} Y_2, W_1),
$$
  

$$
-g_1(\nabla_{Y_1} W_1, Y_2) = -g_1(Y_1, Y_2)g_1(gradh, W_1)
$$

for  $Y_1, Y_2 \in \Gamma(\mathfrak{D}_2)$  and  $W_1 \in \Gamma(\ker F_*)^{\perp} \backslash {\xi}.$ 

Using  $(2.3)$  and  $(2.6)$  in above equation, we get

<span id="page-10-0"></span>
$$
g_1(\nabla_{W_1} \phi Y_1, \phi Y_2) = g_1(\phi Y_1, \phi Y_2) g_1(gradh, W_1).
$$
\n(3.49)

Since F is  $CSI$ – submersion and using (3.[49\)](#page-10-0), we have

<span id="page-10-1"></span>
$$
g_2(F_*(\nabla_{W_1}^F \phi Y_1), F_*(\phi Y_2)) = g_2(F_*(\phi Y_1), F_*(\phi Y_2))g_1(gradh, W_1). \tag{3.50}
$$

From  $(2.17)$  $(2.17)$  in  $(3.50)$  $(3.50)$ , we obtain

$$
g_2(\nabla_{W_1} F_*(\phi Y_1), F_*(\phi Y_2)) = g_2(F_*(\phi Y_1), F_*(\phi Y_2))g_1(gradh, W_1),
$$
\n(3.51)

which implies  $\overline{\nabla}_{W_1} F_*(\phi Y_1) = W_1(h) F_*(\phi Y_1)$  for  $Y_1 \in \Gamma(\mathfrak{D}_2)$  and  $W_1 \in \Gamma(\ker F_*)^{\perp} \setminus {\{\xi\}}$ .

**Theorem 3.4.** Let F be a CSI- submersion with  $r = e^h$  from a Cosymplectic manifold  $(\mathcal{N}_1, \phi, \xi, \eta, g_1)$  onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$ . If  $\mathcal T$  is not equal to zero identically, then the invariant distribution  $\mathfrak{D}_1$  cannot defined a totally geodesic foliation on  $\mathcal{N}_1$ .

**Proof.** For 
$$
Y_1, Y_2 \in \Gamma(\mathfrak{D}_1)
$$
 and  $U_1 \in \Gamma(\mathfrak{D}_2)$ , using (2.3), (2.6), (2.13) and (2.16), we

get

$$
g_1(\nabla_{Y_1} Y_2, U_1) = g_1(\nabla_{Y_1} \phi Y_2, \phi U_1),
$$
  
=  $g_1(\mathcal{T}_{Y_1} \phi Y_2, \phi U_1),$   
=  $-g_1(Y_1, \phi Y_2)g_1(gradh, \phi U_1).$ 

Thus, one can easily obtain the assertion from above equation and the fact that  $gradh \in$  $\phi(\mathfrak{D}_2).$ 

**Theorem 3.5.** The CSI – submersion F with  $r = e^h$  from a Cosymplectic manifold  $(\mathcal{N}_1, \phi, \xi, \phi)$  $(\eta, g_1)$  onto a Riemannian manifold  $(\mathcal{N}_2, g_2)$ . Then the fibers of F are totally geodesic or the anti-invariant distribution  $\mathfrak{D}_2$  is one-dimensional.

**Proof.** The result is quite obvious when we take the fibers of  $F$  are totally geodesic. For second one, since F is a CSI- submersion, then either  $\dim(\mathfrak{D}_2) = 1$  or  $\dim(\mathfrak{D}_2) > 1$ . If  $\dim(\mathcal{D}_2) > 1$ , then we can choose  $U_1, U_2 \in \Gamma(\mathfrak{D}_2)$  such that  $\{U_1, U_2\}$  is orthonormal. From  $(2.13), (3.19)$  $(2.13), (3.19)$  $(2.13), (3.19)$  $(2.13), (3.19)$  and  $(3.20)$  $(3.20)$ , we get

$$
\begin{array}{rcl}\n\mathcal{T}_{U_1}\phi U_2 + \mathcal{H}\nabla_{U_1}\phi U_2 &=& \nabla_{U_1}\phi U_2, \\
\mathcal{T}_{U_1}\phi U_2 + \mathcal{H}\nabla_{U_1}\phi U_2 &=& \mathcal{B}\mathcal{T}_{U_1}U_2 + \mathcal{C}\mathcal{T}_{U_1}U_2 + \psi\mathcal{V}\nabla_{U_1}U_2 + \omega\mathcal{V}\nabla_{U_1}U_2.\n\end{array}
$$

Taking the inner product above equation with  $U_1$ , we obtain

$$
g_1(\mathcal{T}_{U_1}\phi U_2, U_1) = g_1(B\mathcal{T}_{U_1}U_2, U_1) + g_1(\psi V \nabla_{U_1} U_2, U_1).
$$
\n(3.52)

From  $(2.3), (2.6)$  $(2.3), (2.6)$  $(2.3), (2.6)$  $(2.3), (2.6)$  and  $(2.13),$  we have

<span id="page-11-0"></span>
$$
g_1(\mathcal{T}_{U_1}U_1, \phi U_2) = -g_1(\mathcal{T}_{U_1}\phi U_2, U_1) = g_1(\mathcal{T}_{U_1}U_2, \phi U_1).
$$
\n(3.53)

Now, using  $(2.16)$  $(2.16)$  and  $(3.53)$  $(3.53)$ , we get

<span id="page-11-1"></span>
$$
g_1(\mathcal{T}_{U_1}U_1, \phi U_2) = -g_1(gradh, \phi U_2). \tag{3.54}
$$

From equations  $(2.16)$  $(2.16)$  and  $(3.54)$  $(3.54)$ , we obtain

$$
-g_1(gradh, \phi U_2) = g_1(\mathcal{T}_{U_1}U_1, \phi U_2) = -g_1(\mathcal{T}_{U_1} \phi U_2, U_1) = g_1(\mathcal{T}_{U_1}U_2, \phi U_1).
$$
(3.55)

From above equation, we get

$$
g_1(gradh, \phi U_2) = -g_1(\mathcal{T}_{U_1} U_2, \phi U_1),
$$
  
\n $g_1(gradh, \phi U_2) = g_1(U_1, U_2)g_1(gradh, \phi U_1),$   
\n $g_1(gradh, \phi U_2) = 0.$ 

Thus, we get  $\text{grad} h \perp \phi(\mathfrak{D}_2)$ .

Therefore, the dimension of  $\mathfrak{D}_2$  must be one.

#### 4. Example

**Example 4.1.** Taking an Euclidean space  $\mathcal{N}_1$ , given by  $\mathcal{N}_1 = \{(x_1, x_2, y_1, y_2, z) \in \mathbb{R}^5$ :  $(x_1, x_2, y_1, y_2) \neq (0, 0, 0, 0)$  and  $z \neq 0$ . We define the Riemannian metric  $g_1$  on  $\mathcal{N}_1$  defined as  $g_1 = e^{2z} dx_1^2 + e^{2z} dx_2^2 + e^{2z} dy_1^2 + e^{2z} dy_2^2 + dz^2$  and the Cosymplectic structure on  $\phi$  and  $\mathcal{N}_1$ defined as  $\phi(x_1, x_2, y_1, y_2, z) = (y_1, y_2, -x_1, -x_2, z).$ 

Let  $\mathcal{N}_2 = \{(v_1, v_2) \in \mathbb{R}^2\}$  be a Riemannian manifold with Riemannian metric  $g_2$ , given by  $g_2 = e^{2z} dv_1^2 + dv_2^2$ . Define a map  $F : R^5 \to R^2$  by

$$
F(x_1, x_2, y_1, y_2, z) = \left(\frac{x_2 - y_2}{\sqrt{2}}, z\right).
$$

Then, we have

$$
\ker F_* = \langle X_1 = e_1, X_2 = e_2 + e_4, X_3 = e_3 \rangle,
$$
  

$$
\mathfrak{D}_1 = \langle X_1 = e_1, X_3 = e_3 \rangle, \mathfrak{D}_2 = \langle X_2 = e_2 + e_4 \rangle,
$$
  

$$
(\ker F_*)^{\perp} = \langle H_1 = e_2 - e_4, H_2 = e_5 \rangle,
$$

where  $\{e_1 = e^{-z} \frac{\partial}{\partial x}$  $\frac{\partial}{\partial x_1}, e_2 = e^{-z} \frac{\partial}{\partial x}$  $\frac{\partial}{\partial x_2}, e_3 = e^{-z} \frac{\partial}{\partial y}$  $\frac{\partial}{\partial y_1}, e_4 = e^{-z} \frac{\partial}{\partial y}$  $\frac{\partial}{\partial y_2}, e_7 = \frac{\partial}{\partial z} \},\, \{e_1^* = \frac{\partial}{\partial v}$  $\frac{\partial}{\partial v_1}, e_2^* = \frac{\partial}{\partial v}$  $\frac{\partial}{\partial v_2}\}$ are bases on  $T_p\mathcal{N}_1$  and  $T_{F(p)}\mathcal{N}_2$ , respectively, for all  $p \in \mathcal{N}_1$ . By direct computations, we can see that  $F_*(H_1) = \sqrt{2}e^{-z}e_1^*, F_*(H_2) = e_2^*,$  and  $g_1(H_i, H_j) = g_2(F_*H_i, F_*H_j)$  for all  $H_i, H_j \in \Gamma(\ker F_*)^{\perp}, i, j = 1, 2$ . Thus, F is submersion. Moreover, it is easy to see that  $\phi X_1 = -X_3, \phi X_2 = -H_1$  and  $\phi X_3 = X_1$ . Therefore F is a CSI– submersion.

Now, using the Cosymplectic structure, we see that

<span id="page-12-0"></span>
$$
[e_1, e_1] = [e_2, e_2] = [e_3, e_3] = [e_4, e_4] = [e_5, e_5] = 0,
$$
\n
$$
[e_1, e_2] = 0, [e_1, e_3] = 0, [e_1, e_4] = 0, [e_1, e_5] = e_1,
$$
\n
$$
[e_2, e_3] = 0, [e_2, e_4] = 0, [e_2, e_5] = e_2, [e_3, e_4] = 0,
$$
\n
$$
[e_3, e_5] = e_3, [e_4, e_5] = e_4.
$$
\n(4.56)

The Levi-Civita connection  $\nabla$  of the metric  $g_1$  is given by the Koszul's formula which is

<span id="page-12-1"></span>
$$
2g_1(\nabla_X Y, Z)
$$
\n
$$
= Xg_1(Y, Z) + Yg_1(Z, X) - Zg_1(X, Y) + g_1([X, Y], Z) - g_1([Y, Z], X) + g_1([Z, X], Y).
$$
\n
$$
(4.57)
$$

Using equations  $(4.56)$  $(4.56)$  and  $(4.57)$  $(4.57)$ , we obtain

$$
\nabla_{e_1} e_1 = \nabla_{e_2} e_2 = \nabla_{e_3} e_3 = \nabla_{e_4} e_4 = -\frac{\partial}{\partial z},
$$
\n(4.58)  
\n
$$
\nabla_{e_1} e_2 = \nabla_{e_1} e_3 = \nabla_{e_1} e_4 = \nabla_{e_2} e_1 = \nabla_{e_2} e_3 = \nabla_{e_2} e_4 = 0,
$$
\n
$$
\nabla_{e_3} e_1 = \nabla_{e_3} e_2 = \nabla_{e_3} e_4 = \nabla_{e_4} e_1 = \nabla_{e_4} e_2 = \nabla_{e_4} e_3 = 0.
$$
\n
$$
\nabla_{e_1} e_5 = e_1, \nabla_{e_2} e_5 = e_2, \nabla_{e_3} e_5 = e_3, \nabla_{e_4} e_5 = e_4, \nabla_{e_5} e_5 = 0.
$$

Therefore, we have

<span id="page-13-0"></span>
$$
\nabla_{X_1} X_1 = \nabla_{e_1} e_1 = -\frac{\partial}{\partial z}, \nabla_{X_2} X_2 = \nabla_{e_2 + e_4} e_2 + e_4 = -2\frac{\partial}{\partial z},
$$
(4.59)  
\n
$$
\nabla_{X_3} X_3 = \nabla_{e_3} e_3 = -2\frac{\partial}{\partial z}, \nabla_{X_1} X_2 = \nabla_{e_1} e_2 = \nabla_{X_1} X_3 = \nabla_{e_1} e_3 = 0,
$$
  
\n
$$
\nabla_{X_2} X_3 = \nabla_{e_2} e_3 = 0, \nabla_{X_2} X_1 = \nabla_{e_2} e_1 = 0, \nabla_{X_3} X_1 = \nabla_{e_3} e_1 = 0,
$$
  
\n
$$
\nabla_{X_3} X_2 = \nabla_{e_3} e_2 + e_4 = 0.
$$

Thus, we have

<span id="page-13-1"></span>
$$
\mathcal{T}_{V}V = \mathcal{T}_{\lambda_{1}X_{1} + \lambda_{2}X_{2} + \lambda_{3}X_{3}}\lambda_{1}V_{1} + \lambda_{2}V_{2} + \lambda_{3}V_{3}, \lambda_{1}, \lambda_{2}, \lambda_{3} \in R,
$$
  
\n
$$
\mathcal{T}_{V}V = \lambda_{1}^{2}\mathcal{T}_{X_{1}}X_{1} + \lambda_{2}^{2}\mathcal{T}_{X_{2}}X_{2} + \lambda_{3}^{2}\mathcal{T}_{X_{3}}X_{3} + \lambda_{1}\lambda_{2}\mathcal{T}_{X_{1}}X_{2} + \lambda_{1}\lambda_{3}\mathcal{T}_{X_{1}}X_{3} + \lambda_{2}\lambda_{3}\mathcal{T}_{X_{2}}X_{3} + \lambda_{1}\lambda_{2}\mathcal{T}_{X_{2}}X_{1} + \lambda_{1}\lambda_{3}\mathcal{T}_{X_{3}}X_{1} + \lambda_{2}\lambda_{3}\mathcal{T}_{X_{3}}X_{2}.
$$
  
\n(4.60)

Using equations  $(2.12)$  $(2.12)$  and  $(4.59)$  $(4.59)$ , we obtain

<span id="page-13-2"></span>
$$
\mathcal{T}_{X_1} X_1 = -\frac{\partial}{\partial z}, \mathcal{T}_{X_2} X_2 = -2\frac{\partial}{\partial z}, \mathcal{T}_{X_3} X_3 = -\frac{\partial}{\partial z}, \n\mathcal{T}_{X_1} X_2 = 0, \mathcal{T}_{X_1} X_3 = 0, \mathcal{T}_{X_2} X_3 = 0, \mathcal{T}_{X_2} X_1 = 0, \n\mathcal{T}_{X_2} X_3 = 0, \mathcal{T}_{X_3} X_1 = 0.
$$
\n(4.61)

Now using equations (4.[60\)](#page-13-1) and (4.[61\)](#page-13-2), we get

$$
\mathcal{T}_V V = -(\lambda_1^2 + 2\lambda_2^2 + \lambda_3^2) \frac{\partial}{\partial z}.
$$
\n(4.62)

Since  $X = \lambda_1 X_1 + \lambda_2 X_2 + \lambda_3 X_3$ , so  $g_1(\lambda_1 V_1 + \lambda_2 V_2 + \lambda_3 V_3, \lambda_1 V_1 + \lambda_2 V_2 + \lambda_3 V_3) = \lambda_1^2 + \lambda_2^2 V_1 + \lambda_3^2 V_2 + \lambda_4^2 V_3$  $2\lambda_2^2 + \lambda_3^2$ . For a smooth function h on  $R^5$ , the  $\nabla h$  w. r. t. the metric  $g_1$  is given by  $\nabla h$  =  $e^{-2z}\frac{\partial h}{\partial x}$  $\partial x_1$ ∂  $\frac{\partial}{\partial x_1} + e^{-2z} \frac{\partial h}{\partial x_2}$  $\partial x_2$ ∂  $\frac{\partial}{\partial x_2} + e^{-2z} \frac{\partial h}{\partial y_1}$  $\partial y_1$ ∂  $\frac{\partial}{\partial y_1} + e^{-2z} \frac{\partial h}{\partial y_2}$  $\partial y_2$ ∂  $\frac{\partial}{\partial y_2} + \frac{\partial h}{\partial z}$ ∂z  $\frac{\partial}{\partial z}$ . Hence  $\nabla h = \frac{\partial}{\partial z}$  for the function h = z. Then one can easily find that  $\mathcal{T}_V V = -g_1(V, V) \nabla h$ , thus by Theorem 3.1, the map F is a CSI− submersion from Cosymplectic manifold onto Riemannian manifold.

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Department of Mathematics, Shri Jai Narain Post Graduate College, Lucknow (U.P.)-India

Department of Mathematics, Dr. S.K.S. Womens College Motihari, B.R.Ambedkar Bihar University Muzaffarpur-India

Department of Mathematics, Shri Jai Narain Post Graduate College, Lucknow (U.P.)-India