



## FRENET CURVES IN 3-DIMENSIONAL LORENTZIAN CONCIRCULAR STRUCTURE MANIFOLDS

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**ABSTRACT.** In this paper, we give some characterizations of Frenet curves in 3-dimensional Lorentzian concircular structure manifolds( $(LCS)_3$  manifolds). We define Frenet equations and the Frenet elements of these curves. We also obtain the curvatures of non-geodesic Frenet curves on  $(LCS)_3$  manifolds. Finally we give some theorems, corollaries and examples for these curves.

**Keywords:** Lorentzian manifold, Concircular structure, Frenet curve

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### 1. INTRODUCTION

The differential geometry of curves in manifolds investigated by several authors. Especially the curves in contact and para-contact manifolds drew attention and studied by the authors. B. Olszak[17], derived the conditions for an a.c.m structure on M to be normal and point out some of their consequences. B. Olszak completely characterized the local nature of normal a.c.m. structures on M by giving suitable examples. Moreover B. Olszak gave some restrictions on the scalar curvature in contact metric manifolds which are conformally flat or of constant  $\phi$ -sectional curvature in[16].

J. Welyczko[22], generalized some of the results for Legendre curves in three dimensional normal a.c.m. manifolds, especially, quasi-Sasakian manifolds. J. Welyczko [23], studied

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the curvatures of slant Frenet curves in three-dimensional normal almost paracontact metric manifolds.

B. E. Acet and S. Y. Perktaş [1] obtained the curvatures of Legendre curves in 3-dimensional  $(\varepsilon, \delta)$  trans-Sasakian manifolds. Ji-Eun Lee, defined Lorentzian cross product in a three-dimensional almost contact Lorentzian manifold and proved that  $\frac{\kappa}{\tau-1} = \text{cons.}$  along a Frenet slant curve in a Sasakian Lorentzian three-manifold. Furthermore, Ji-Eun Lee proved that  $\gamma$  is a slant curve if and only if  $M$  is Sasakian for a contact magnetic curve  $\gamma$  in contact Lorentzian 3-manifold  $M$  in[12]. Ji-Eun Lee, also gave some characterizations for the generalized Tanaka-Webster connection in a contact Lorentzian manifold in[13].

A. Yıldırım[25] obtained the Frenet apparatus for Frenet curves on three dimensional normal almost contact manifolds and characterized some results for these curves.

U.C.De and K.De[10] studied Lorentzian Trans-Sasakian and conformally flat Lorentzian Trans-Sasakian manifolds.

The LCS manifolds was introduced by [19] with an example. A. A. Shaikh[20] studied various types of  $(LCS)_n$ -manifolds and proved that in such a manifold the Ricci operator commutes with the structure tensor  $\varphi$ .

In this framework, the paper is organized in the following way. Section 2 with two subsections, we give basic definitions of a  $(LCS)_n$ -manifolds manifold. In the second subsection we give the Frenet-Serret equations of a curve in  $(LCS)_3$  manifold. We give finally the Frenet elements of a Frenet curve in  $(LCS)_3$  manifold and give theorems, corollaries and examples for these curves in the third and fourth sections.

## 2. PRELIMINARIES

**2.1. Lorentzian Conircular Structure Manifolds.** A Lorentzian manifold of dimension  $n$  is a doublet  $(\bar{N}, \bar{g})$ , where  $\bar{N}$  is a smooth connected para-compact Hausdorff manifold of dimension  $n$  and  $\bar{g}$  is a Lorentzian metric, that is,  $\bar{N}$  admits a smooth symmetric tensor field  $g$  of type  $(0, 2)$  such that for each point  $p \in \bar{N}$  the tensor  $\bar{g}_p : T_p\bar{N} \times T_p\bar{N} \rightarrow R$  is a non degenerate inner product of signature  $(-, +, \dots, +)$ , where  $T_p\bar{N}$  denotes the tangent space of  $\bar{N}$  at  $p$  and  $R$  is the real number space. A non zero vector field  $V \in T_p\bar{N}$  is called spacelike (resp. non-spacelike, null and timelike) if it satisfies  $\bar{g}_p(V, V) > 0$  (resp.,  $\leq 0, =, < 0$ ).[15]

**Definition 2.1.** In a Lorentzian manifold  $(\bar{N}, \bar{g})$  a vector field  $w$  is defined by

$$\bar{g}(U, \rho) = A(U) \tag{2.1}$$

for any  $U \in \chi(\bar{N})$  is said to be a concircular vector field if

$$(\nabla_U A)(V) = \alpha \{ \bar{g}(U, V) + w(U)w(V) \}, \tag{2.2}$$

where  $\alpha$  is a non-zero scalar and  $w$  is a closed 1-form.[24]

If a Lorentzian manifold  $\bar{N}$  admits a unit timelike concircular vector field  $\xi$ , called generator of the manifold, then we have

$$\bar{g}(\xi, \xi) = -1. \tag{2.3}$$

Since  $\xi$  is the unit concircular vector field on  $\bar{N}$ , there exists a non-zero 1-form  $\eta$  such that

$$\bar{g}(U, \xi) = \eta(U), \tag{2.4}$$

which satisfies the following equation

$$(\nabla_U \eta)(V) = \alpha \{ \bar{g}(U, V) + \eta(U)\eta(V) \}, \quad (\alpha \neq 0) \tag{2.5}$$

for all vector fields  $U$  and  $V$ , where  $\nabla$  gives the covariant differentiation with respect to the Lorentzian metric  $\bar{g}$  and  $\alpha$  is a non-zero scalar function satisfies

$$(\nabla_U \alpha) = U\alpha = d\alpha(U) = \rho\eta(U), \tag{2.6}$$

where  $\rho$  is a certain scalar function defined by  $\rho = -(\xi\alpha)$ . If we take

$$\varphi U = \frac{1}{\alpha} \nabla_U \xi, \tag{2.7}$$

then with the help of (2.3), (2.4) and (2.6), we can find

$$\varphi U = U + \eta(U)\xi, \tag{2.8}$$

which shows that  $\varphi$  is a tensor field of type (1,1), called the structure tensor of the manifold  $\bar{N}$ . Hence the Lorentzian manifold  $\bar{N}$  of class  $C^\infty$  equipped with a unit timelike concircular vector field  $\xi$ , its associated 1-form  $\eta$  and (1,1) tensor field  $\varphi$  is said to be a Lorentzian concircular structure manifold (i.e.  $(LCS)_n$  manifold)[19]. Moreover, if  $\alpha = 1$ , then we have the LP-Sasakian structure of Matsumoto[14]. So we can say the generalization of LP-Sasakian manifold gives us the  $(LCS)_n$  manifold. It is noteworthy to mention that LCS-manifold is invariant under a conformal change whereas LP-Sasakian structure is not so[18]. In  $(LCS)_3$  manifolds, the following relations hold[19]

$$\varphi^2 U = U + \eta(U)\xi, \quad \eta(\xi) = -1, \tag{2.9}$$

$$\varphi(\xi) = 0, \quad \eta(\varphi U) = 0,$$

and

$$\bar{g}(\varphi U, \varphi V) = \bar{g}(U, V) + \eta(U)\eta(V). \quad (2.10)$$

**2.2. Frenet Curves.** Let  $\zeta : I \rightarrow \bar{N}$  be a unit speed curve in  $(LCS)_3$  manifold  $\bar{N}$  such that  $\zeta'$  satisfies  $\bar{g}(\zeta', \zeta') = \varepsilon_1 = \mp 1$ . The constant  $\varepsilon_1$  is called the casual character of  $\zeta$ . The constants  $\varepsilon_2$  and  $\varepsilon_3$  defined by  $\bar{g}(n, n) = \varepsilon_2$  and  $\bar{g}(b, b) = \varepsilon_3$  and called the second casual character and third casual character of  $\zeta$ , respectively. Thus we  $\varepsilon_1\varepsilon_2 = -\varepsilon_3$ .

A unit speed curve  $\zeta$  is said to be a spacelike or timelike if its casual character is 1 or -1, respectively. A unit speed curve  $\zeta$  is said to be a Frenet curve if  $\bar{g}(\zeta', \zeta') \neq 0$ . A Frenet curve  $\zeta$  admits an orthonormal frame field  $\{t = \zeta', n, b\}$  along  $\zeta$ . Then the Frenet-Serret equations given as follows:

$$\begin{aligned} \nabla_{\zeta'} t &= \varepsilon_2 \kappa n \\ \nabla_{\zeta'} n &= -\varepsilon_1 \kappa t - \varepsilon_3 \tau b \\ \nabla_{\zeta'} b &= \varepsilon_2 \tau n \end{aligned} \quad (2.11)$$

where  $\kappa = |\nabla_{\zeta'} \zeta'|$  is the geodesic curvature of  $\zeta$  and  $\tau$  is geodesic torsion [12]. The vector fields  $t$ ,  $n$  and  $b$  are called the tangent vector field, the principal normal vector field and the binormal vector field of  $\zeta$ , respectively.

If the geodesic curvature of the curve  $\zeta$  vanishes, then the curve is called a geodesic curve. If  $\kappa = \text{cons.}$  and  $\tau = 0$ , then the curve is called a pseudo-circle and pseudo-helix if the geodesic curvature and torsion are constant.

A curve in a three dimensional Lorentzian manifold is a slant curve if the tangent vector field of the curve has constant angle with the Reeb vector field, i.e.  $\eta(\zeta') = -\bar{g}(\zeta', \xi) = \cos\theta = \text{constant}$ . If  $\eta(\zeta') = -\bar{g}(\zeta', \xi) = 0$ , then the curve  $\zeta$  is called a Legendre curve [12].

### 3. MAIN RESULTS

In this section we consider a  $(LCS)_3$  manifold  $\bar{N}$ . Let  $\zeta : I \rightarrow \bar{N}$  be a Frenet curve with the geodesic curvature  $\kappa \neq 0$ , given with the arc-parameter  $s$  and  $\bar{\nabla}$  be the Levi-Civita connection on  $\bar{N}$ . From the basis  $(\zeta', \varphi\zeta', \xi)$  we obtain an orthonormal basis  $\{e_1, e_2, e_3\}$  which

satisfy the equations

$$\begin{aligned} e_1 &= \zeta', \\ e_2 &= \frac{\varepsilon_2 \varphi \zeta'}{\sqrt{\varepsilon_1 + \rho^2}}, \\ e_3 &= \varepsilon_2 \frac{\varepsilon_1 \xi - \rho \zeta'}{\sqrt{\varepsilon_1 + \rho^2}} \end{aligned} \tag{3.12}$$

where

$$\eta(\zeta') = \bar{g}(\zeta', \xi) = \rho. \tag{3.13}$$

Then if we write the covariant differentiation of  $\zeta'$  as

$$\bar{\nabla}_{\zeta'} e_1 = \nu e_2 + \mu e_3 \tag{3.14}$$

such that

$$\nu = \bar{g}(\bar{\nabla}_{\zeta'} e_1, e_2) \tag{3.15}$$

is a certain function. Moreover we obtain  $\nu$  by

$$\mu = \bar{g}(\bar{\nabla}_{\zeta'} e_1, e_3) = \varepsilon_2 \left( \frac{\rho'}{\sqrt{\varepsilon_1 + \rho^2}} - \varepsilon_1 \alpha \sqrt{\varepsilon_1 + \rho^2} \right), \tag{3.16}$$

where  $\rho'(s) = \frac{d\rho(\zeta(s))}{ds}$ . Then we find

$$\bar{\nabla}_{\zeta'} e_2 = -\nu e_1 + \left( \varepsilon_3 \alpha + \frac{\varepsilon_1 \rho \nu}{\sqrt{\varepsilon_1 + \rho^2}} \right) e_3 \tag{3.17}$$

and

$$\bar{\nabla}_{\zeta'} e_3 = -\mu e_1 - \left( \varepsilon_3 \alpha + \frac{\varepsilon_1 \rho \nu}{\sqrt{\varepsilon_1 + \rho^2}} \right) e_2. \tag{3.18}$$

The fundamental forms of the tangent vector  $\zeta'$  on the basis of the equation (3.12) is

$$[\omega_{ij}(\zeta')] = \begin{pmatrix} 0 & \nu & \mu \\ -\nu & 0 & \varepsilon_3 \alpha + \frac{\varepsilon_1 \rho \nu}{\sqrt{\varepsilon_1 + \rho^2}} \\ -\mu & -\varepsilon_3 \alpha - \frac{\varepsilon_1 \rho \nu}{\sqrt{\varepsilon_1 + \rho^2}} & 0 \end{pmatrix} \tag{3.19}$$

and the Darboux vector connected to the vector  $\zeta'$  is

$$\omega(\zeta') = \left( \varepsilon_3 \alpha + \frac{\varepsilon_1 \rho \nu}{\sqrt{\varepsilon_1 + \rho^2}} \right) e_1 - \mu e_2 + \nu e_3. \tag{3.20}$$

So we can write

$$\bar{\nabla}_{\zeta'} e_i = \omega(\zeta') \wedge \varepsilon_i e_i \quad (1 \leq i \leq 3). \tag{3.21}$$

Thus, for any vector field  $Z = \sum_{i=1}^3 \theta^i e_i \in \chi(\bar{N})$  strictly dependent on the curve  $\zeta$  on  $\bar{N}$  and we have the following equation

$$\bar{\nabla}_{\zeta'} Z = \omega(\zeta') \wedge Z + \sum_{i=1}^3 \varepsilon_i e_i [\theta^i] e_i. \quad (3.22)$$

**3.1. Frenet Elements of  $\zeta$ .** Let a curve  $\zeta : I \rightarrow \bar{N}$  be a Frenet curve with the geodesic curvature  $\kappa \neq 0$ , given with the arc parameter  $s$  and the elements  $\{t, n, b, \kappa, \tau\}$ . The Frenet elements of the curve  $\zeta$  can be calculated as above:

If we consider the equation (3.14), then we get

$$\varepsilon_2 \kappa n = \bar{\nabla}_{\zeta'} e_1 = \nu e_2 + \mu e_3. \quad (3.23)$$

If we consider (3.16) and (3.23) we find

$$\kappa = \sqrt{\nu^2 + \left( \frac{\rho'}{\sqrt{\varepsilon_1 + \rho^2}} - \varepsilon_1 \alpha \sqrt{\varepsilon_1 + \rho^2} \right)^2}. \quad (3.24)$$

On the other hand

$$\begin{aligned} \bar{\nabla}_{\zeta'} n &= \left( \frac{\nu}{\varepsilon_2 \kappa} \right)' e_2 + \frac{\nu}{\varepsilon_2 \kappa} \nabla_{\zeta'} e_2 + \left( \frac{\mu}{\varepsilon_2 \kappa} \right)' e_3 + \frac{\mu}{\varepsilon_2 \kappa} \nabla_{\zeta'} e_3 \\ &= -\varepsilon_1 \kappa t - \varepsilon_3 \tau B. \end{aligned} \quad (3.25)$$

By means of the equation (3.17) and (3.18) we find

$$\begin{aligned} -\varepsilon_3 \tau B &= \left[ \left( \frac{\nu}{\varepsilon_2 \kappa} \right)' - \frac{\mu}{\varepsilon_2 \kappa} \left( \varepsilon_3 \alpha + \frac{\varepsilon_1 \rho \nu}{\sqrt{\varepsilon_1 + \rho^2}} \right) \right] e_2 \\ &+ \left[ \left( \frac{\mu}{\varepsilon_2 \kappa} \right)' + \frac{\nu}{\varepsilon_2 \kappa} \left( \varepsilon_3 \alpha + \frac{\varepsilon_1 \rho \nu}{\sqrt{\varepsilon_1 + \rho^2}} \right) \right] e_3. \end{aligned} \quad (3.26)$$

By a direct computation we find following

$$\left[ \left( \frac{\nu}{\varepsilon_2 \kappa} \right)' \right]^2 + \left[ \left( \frac{\mu}{\varepsilon_2 \kappa} \right)' \right]^2 = \left[ - \left( \frac{\nu}{\varepsilon_2 \kappa} \right)' \frac{\mu}{\varepsilon_2 \kappa} + \frac{\nu}{\varepsilon_2 \kappa} \left( \frac{\mu}{\varepsilon_2 \kappa} \right)' \right]^2. \quad (3.27)$$

Taking the norm of the last equation by using (3.26) and if we consider the equations (3.16) and (3.27) in (3.26) we obtain

$$\tau = \left| \varepsilon_3 \alpha + \frac{\varepsilon_1 \rho \nu}{\sqrt{\varepsilon_1 + \rho^2}} - \sqrt{\left[ \left( \frac{\nu}{\varepsilon_2 \kappa} \right)' \right]^2 + \left[ \left( \frac{\varepsilon_2 \left( \frac{\rho'}{\sqrt{\varepsilon_1 + \rho^2}} - \varepsilon_1 \alpha \sqrt{\varepsilon_1 + \rho^2} \right)}{\kappa} \right)' \right]^2} \right|. \quad (3.28)$$

Moreover we can write the Frenet vector fields of  $\zeta$  as in the following theorem

**Theorem 3.1.** *Let  $\bar{N}$  be a  $(LCS)_3$  manifold and  $\zeta$  be a Frenet curve on  $\bar{N}$ . The Frenet vector fields  $t$ ,  $n$  and  $b$  are in the form of*

$$\begin{aligned} t &= \zeta' = e_1, \\ n &= \frac{\nu}{\varepsilon_2\kappa}e_2 + \frac{\mu}{\varepsilon_2\kappa}e_3, \\ b &= -\frac{1}{\varepsilon_3\tau} \left[ \left( \frac{\nu}{\varepsilon_2\kappa} \right)' - \frac{\mu}{\varepsilon_2\kappa} \left( \varepsilon_3\alpha + \frac{\varepsilon_1\rho\nu}{\sqrt{\varepsilon_1 + \rho^2}} \right) \right] e_2 \\ &\quad - \frac{1}{\varepsilon_3\tau} \left[ \left( \frac{\mu}{\varepsilon_2\kappa} \right)' + \frac{\nu}{\varepsilon_2\kappa} \left( \varepsilon_3\alpha + \frac{\varepsilon_1\rho\nu}{\sqrt{\varepsilon_1 + \rho^2}} \right) \right] e_3. \end{aligned} \tag{3.29}$$

Note that

$$\begin{aligned} \xi &= \varepsilon_1\rho t - \frac{\mu\sqrt{\varepsilon_1 + \rho^2}}{\kappa}n \\ &\quad - \frac{\sqrt{\varepsilon_1 + \rho^2}}{\varepsilon_3\tau} \left[ \left( \frac{\mu}{\varepsilon_2\kappa} \right)' + \frac{\nu}{\varepsilon_2\kappa} \left( \varepsilon_3\alpha + \frac{\varepsilon_1\rho\nu}{\sqrt{\varepsilon_1 + \rho^2}} \right) \right] b. \end{aligned} \tag{3.30}$$

Let  $\zeta$  be a non-geodesic Frenet curve given with the arc-parameter  $s$  in  $(LCS)_3$  manifold  $\bar{N}$ . So one can state the above theorems.

**Theorem 3.2.** *Let  $\bar{N}$  be a  $(LCS)_3$  manifold and  $\zeta$  be a Frenet curve on  $\bar{N}$ .  $\zeta$  is a slant curve ( $\rho = \eta(\zeta') = \cos\theta = \text{const.}$ ) on  $\bar{N}$  if and only if the Frenet elements  $\{t, n, b, \kappa, \tau\}$  of  $\zeta$  are as follows*

$$\begin{aligned} t &= e_1 = \zeta', \\ n &= e_2 = \frac{\varepsilon_2\varphi\zeta'}{\sqrt{\varepsilon_1 + \cos^2\theta}}, \\ b &= e_3 = \varepsilon_2 \frac{\varepsilon_1\xi - \cos\theta\zeta'}{\sqrt{\varepsilon_1 + \cos^2\theta}}, \\ \kappa &= \sqrt{\nu^2 + \alpha^2(\varepsilon_1 + \cos^2\theta)}, \\ \tau &= \left| \varepsilon_3\alpha + \frac{\varepsilon_1\cos\theta\nu}{\sqrt{\varepsilon_1 + \cos^2\theta}} - \sqrt{\left[ \left( \frac{\nu}{\varepsilon_2\kappa} \right) \right]^2 + \left[ \left( \frac{\alpha\sqrt{\varepsilon_1 + \cos^2\theta}}{\kappa} \right) \right]^2} \right|. \end{aligned} \tag{3.31}$$

**Proof.** Let the curve  $\zeta$  be a slant curve in  $(LCS)_3$  manifold  $\bar{N}$ . If we take account the condition  $\rho = \eta(\zeta') = \cos\theta = \text{constant}$  in the equations (3.12), (3.24) and (3.28) we find (3.31). If the equations in (3.31) hold, from the definition of slant curves it is obvious that the curve  $\zeta$  is a slant curve.

**Corollary 3.1.** *Let  $\bar{N}$  be a  $(LCS)_3$  manifold and  $\zeta$  be a slant curve on  $\bar{N}$ . If the geodesic curvature  $\kappa$  of the curve  $\zeta$  is non-zero constant, then the geodesic torsion of  $\zeta$  is  $\tau = \left| \left( \varepsilon_3 \alpha + \varepsilon_1 \frac{\cos \theta \nu}{\sqrt{\varepsilon_1 + \cos^2 \theta}} \right) \right|$  and  $\zeta$  is a pseudo-helix on  $\bar{N}$ .*

**Corollary 3.2.** *Let  $\bar{N}$  be a  $(LCS)_3$  manifold and  $\zeta$  be a slant curve on  $\bar{N}$ . If the geodesic curvature  $\kappa$  of the curve  $\zeta$  is not constant and the geodesic torsion of  $\zeta$  is  $\tau = 0$  then  $\zeta$  is a plane curve on  $\bar{N}$  and function  $\nu$  satisfies the equation*

$$\nu = \int (c_1 + c_2 \nu) \kappa^2 ds, \quad (3.32)$$

where  $c_1 = \frac{\varepsilon_3}{\sqrt{\varepsilon_1 + \cos^2 \theta}}$  and  $c_2 = \frac{\varepsilon_1 \cos \theta}{\alpha(\varepsilon_1 + \cos^2 \theta)}$ .

**Theorem 3.3.** *Let  $\bar{N}$  be a  $(LCS)_3$  manifold and  $\zeta$  is a Frenet curve on  $\bar{N}$ .  $\zeta$  is a spacelike Legendre curve ( $\rho = \eta(\zeta') = 0$ ) in this manifold if and only if the Frenet elements  $\{t, n, b, \kappa, \tau\}$  of  $\zeta$  are as follows*

$$\begin{aligned} t &= e_1 = \zeta', \\ n &= e_2 = \varepsilon_2 \varphi \zeta', \\ b &= e_3 = -\varepsilon_3 \xi, \\ \kappa &= \sqrt{\nu^2 + \alpha^2}, \\ \tau &= \left| \varepsilon_3 \alpha - \sqrt{\left[ \left( \frac{\nu}{\varepsilon_2 \kappa} \right)' \right]^2 + \alpha^2 \left[ \frac{\kappa'}{\kappa^2} \right]^2} \right|. \end{aligned} \quad (3.33)$$

**Proof.** Let the curve  $\zeta$  be a Legendre curve in  $(LCS)_3$  manifold  $\bar{N}$ . If we take account the condition  $\rho = \eta(\zeta') = 0$  in the equations (3.12), (3.24) and (3.28) we find (3.33). If the equations in (3.33) hold, from the definition of Legendre curves it is obvious that the curve  $\zeta$  is a Legendre curve on  $\bar{N}$ .

**Corollary 3.3.** *Let the curve  $\zeta$  is a Legendre curve in  $(LCS)_3$  manifold  $\bar{N}$ . If the geodesic curvature  $\kappa$  of the curve  $\zeta$  is non-zero constant, then the geodesic torsion of  $\zeta$  is  $\tau = 0$  and  $\zeta$  is a plane curve on  $\bar{N}$ .*

#### 4. EXAMPLES

Let  $\bar{N}$  be the 3-dimensional manifold given

$$\bar{N} = \{(x, y, z) \in \mathfrak{R}^3, z \neq 0\}, \quad (4.34)$$



where  $(x,y,z)$  denote the standart co-ordinates in  $\mathfrak{R}^3$ . Then

$$E_1 = e^z \left( x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y} \right), \quad E_2 = e^z \frac{\partial}{\partial y}, \quad E_3 = \frac{\partial}{\partial z} \tag{4.35}$$

are linearly independent of each point of  $\bar{N}$ . Let  $g$  be the Lorentzian metric tensor defined by

$$\begin{aligned} \bar{g}(E_1, E_1) = \bar{g}(E_2, E_2) = 1, \quad \bar{g}(E_3, E_3) = -1, \\ \bar{g}(E_i, E_j) = 0, \quad i \neq j, \end{aligned} \tag{4.36}$$

for  $i, j = 1, 2, 3$ [2]. Let  $\eta$  be the 1-form defined by  $\eta(Z) = \bar{g}(Z, E_3)$  for any  $Z \in \Gamma(T\bar{N})$ . Let  $\varphi$  be the (1,1)-tensor field defined by

$$\varphi E_1 = E_1, \quad \varphi E_2 = E_2, \quad \varphi E_3 = 0. \tag{4.37}$$

Then using the condition of the linearity of  $\varphi$  and  $\bar{g}$ , we obtain  $\eta(E_3) = -1$ ,

$$\begin{aligned} \varphi^2 Z = Z + \eta(Z)E_3, \\ \bar{g}(\varphi Z, \varphi W) = \bar{g}(Z, W) + \eta(Z)\eta(W), \end{aligned} \tag{4.38}$$

for all  $Z, W \in \Gamma(T\bar{N})$ . Thus for  $\xi = E_3$ ,  $(\varphi, \xi, \eta, \bar{g})$  defines a Lorentzian paracontact structure on  $\bar{N}$ .

Now, let  $\nabla$  be the Levi-Civita connection with respect to the Lorentzian metric  $\bar{g}$ . Then we obtain

$$[E_1, E_2] = -e^z E_2, \quad [E_1, E_3] = -E_1, \quad [E_2, E_3] = -E_2. \tag{4.39}$$

If we use the Koszul formulae for the Lorentzian metric tensor  $\bar{g}$ , we can easily calculate the covariant derivations as follows:

$$\begin{aligned} \nabla_{E_1} E_1 = -E_3, \quad \nabla_{E_2} E_1 = e^z E_2, \quad \nabla_{E_1} E_3 = -E_1, \\ \nabla_{E_2} E_3 = -E_2, \quad \nabla_{E_2} E_2 = -e^z E_1 - E_3, \\ \nabla_{E_1} E_2 = \nabla_{E_3} E_1 = \nabla_{E_3} E_2 = \nabla_{E_3} E_3 = 0. \end{aligned} \tag{4.40}$$

From the about representatations, one can easily see that  $(\varphi, \xi, \eta, \bar{g})$  is a  $(LCS)_3$  structure on  $\bar{N}$ , that is,  $\bar{N}$  is an  $(LCS)_3$ -manifold with  $\alpha = -1$  and  $\rho = 0$ .

**Example 4.1.** Let  $\beta$  be a spacelike Legendre curve in the  $(LCS)_3$  manifold  $\bar{N}$  and defined as

$$\begin{aligned}\beta: I &\rightarrow \bar{N} \\ s &\rightarrow \beta(s) = (s^2, s^2, \ln 2),\end{aligned}$$

where the curve  $\beta$  parametrized by the arc length parameter  $t$ . If we differentiate  $\beta(t)$  and consider (3.12) we find

$$e_1 = \beta'(t), \quad (4.41)$$

$$e_2 = \frac{1}{\sqrt{2}}E_1 + \frac{1}{\sqrt{2}}E_2, \quad (4.42)$$

$$e_3 = \varepsilon_2 E_3. \quad (4.43)$$

If we consider the equations (3.13), (3.14), (3.16), (3.24) and (3.28) we can write

$$\rho = 0, \quad \mu = -\varepsilon_2 \alpha, \quad \nu = -\frac{1}{\sqrt{2}}, \quad (4.44)$$

$$\kappa = \sqrt{\alpha^2 + \frac{1}{2}} = \sqrt{\frac{3}{2}}, \quad \tau = |\alpha| = 1.$$

From the above equations we see that the curve  $\beta$  is a Legendre helix curve in  $\bar{N}$ .

**Example 4.2.** Let  $v$  be a spacelike Legendre curve in the  $(LCS)_3$  manifold  $\bar{N}$  and defined as

$$\begin{aligned}v: I &\rightarrow \bar{N} \\ s &\rightarrow v(s) = (\cos s, \sin s, 1),\end{aligned}$$

where the curve  $v$  parametrized by the arc length parameter  $t$ . If we differentiate  $v(t)$  and consider (3.12) we find

$$e_1 = v'(t), \quad (4.45)$$

$$e_2 = \varepsilon_2 \left( -\sin\left(\frac{t}{e}\right)E_1 + \cos\left(\frac{t}{e}\right)E_2 \right), \quad (4.46)$$

$$e_3 = \varepsilon_2 \partial_3. \quad (4.47)$$

If we consider the equations (3.13), (3.14), (3.16), (3.24) and (3.28) we can write

$$\rho = 0, \quad \mu = -\varepsilon_2 \alpha, \quad \nu = 0, \quad (4.48)$$

$$\kappa = \tau = |\alpha|.$$

So, the curve  $v(s)$  is a Legendre helix curve in  $\bar{N}$ .

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